

The following are derivations of the output loads available in FAST for a 2-bladed turbine configuration. The loads for a 3-bladed turbine are very similar. Note that some of the loads are given multiple names in order to support variation among the user's preferences.

Along with most of the loads are associated partial loads. These partial loads will be used at the end of this document to redevelop portions of the equations of motion to speed up the computations. The definition of these partial loads is as follows:

Let:

$$\mathbf{F}_{Source}^{X_i}(\ddot{q}, \dot{q}, q, t) = \left(\sum_{r=1}^{22} \mathbf{F}_{Source_r}^{X_i}(q, t) \ddot{q}_r \right) + \mathbf{F}_{Source_t}^{X_i}(\dot{q}, q, t)$$

where $\mathbf{F}_{Source_r}^{X_i}$ are the partial forces and $\mathbf{F}_{Source_t}^{X_i}$ is all components of $\mathbf{F}_{Source}^{X_i}$ that are not of this form.

Similarly, let:

$$\mathbf{M}_{Source}^{N_i@X_i}(\ddot{q}, \dot{q}, q, t) = \left(\sum_{r=1}^{22} \mathbf{M}_{Source_r}^{N_i@X_i}(q, t) \ddot{q}_r \right) + \mathbf{M}_{Source_t}^{N_i@X_i}(\dot{q}, q, t)$$

where $\mathbf{M}_{Source_r}^{N_i@X_i}$ are the partial moments and $\mathbf{M}_{Source_t}^{N_i@X_i}$ is all components of $\mathbf{M}_{Source}^{N_i@X_i}$ that are not of this form.

To find the loads characterizing the constraint forces between two bodies, say A and B, all that is needed is to remove body B from the equations of motion and determine what equivalent load applied on A would give the same effect that body B had on A originally.

Blade 1 Root Loads:

There are 10 output loads at the root of blade 1. 5 of them are the 3 components of the root force $\mathbf{F}_{Bl}^{SI}(0)$ (2 components are expressed in both the coned and blade reference frames). The other 5 are the 3 components of the root bending moments, $\mathbf{M}_{Bl}^H(0)$ (again, 2 components are expressed in both the coned and blade reference frames). If blade 1 is to be removed from the turbine, loads $\mathbf{F}_{Bl}^{SI}(0)$ and $\mathbf{M}_{Bl}^H(0)$ applied to the hub at the blade 1 root ($r = 0$) must give the equivalent effect of blade 1 in the resulting equations of motion. The new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Bl} = {}^E \mathbf{v}_r^C \cdot \mathbf{F}_{Bl}^C + {}^E \boldsymbol{\omega}_r^H \cdot \mathbf{M}_{Bl}^H \quad (r = 1, 2, \dots, 22)$$

where the equivalent loads acting at the hub's center of mass (point C) are related to $\mathbf{F}_{Bl}^{SI}(0)$ and $\mathbf{M}_{Bl}^H(0)$ because the hub is rigid as follows:

$$\mathbf{F}_{Bl}^C = \mathbf{F}_{Bl}^{SI}(0) \quad \text{and} \quad \mathbf{M}_{Bl}^H = \mathbf{M}_{Bl}^H(0) + \mathbf{r}^{CSI}(0) \times \mathbf{F}_{Bl}^{SI}(0) \quad \text{or} \quad \mathbf{M}_{Bl}^H = \mathbf{M}_{Bl}^H(0) + [\mathbf{r}^{QSI}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{Bl}^{SI}(0)$$

But since ${}^E \mathbf{v}_r^C = {}^E \mathbf{v}_r^Q + {}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QC}$, this generalized active force can be expanded to:

$$F_r|_{Bl} = ({}^E \mathbf{v}_r^Q + {}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QC}) \cdot \mathbf{F}_{Bl}^{SI}(0) + {}^E \boldsymbol{\omega}_r^H \cdot \{ \mathbf{M}_{Bl}^H(0) + [\mathbf{r}^{QSI}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{Bl}^{SI}(0) \} \quad (r = 1, 2, \dots, 22)$$

Now applying the cyclic permutation law of the scalar triple product:

$$F_r|_{Bl} = {}^E \mathbf{v}_r^Q \cdot \mathbf{F}_{Bl}^{SI}(0) + {}^E \boldsymbol{\omega}_r^H \cdot \{ \mathbf{r}^{QC} \times \mathbf{F}_{Bl}^{SI}(0) \} + {}^E \boldsymbol{\omega}_r^H \cdot \{ \mathbf{M}_{Bl}^H(0) + [\mathbf{r}^{QSI}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{Bl}^{SI}(0) \} \quad (r = 1, 2, \dots, 22)$$

which simplifies to:

$$F_r|_{Bl} = {}^E \mathbf{v}_r^Q \cdot \mathbf{F}_{Bl}^{SI}(0) + {}^E \boldsymbol{\omega}_r^H \cdot [\mathbf{M}_{Bl}^H(0) + \mathbf{r}^{QSI}(0) \times \mathbf{F}_{Bl}^{SI}(0)] \quad (r = 1, 2, \dots, 22)$$

[This can also be simplified to $F_r|_{Bl} = {}^E \mathbf{v}_r^{SI}(0) \cdot \mathbf{F}_{Bl}^{SI}(0) + {}^E \boldsymbol{\omega}_r^H \cdot \mathbf{M}_{Bl}^H(0)$ ($r = 1, 2, \dots, 22$), which will be used later in the ensuing analysis.]

This generalized active force must produce the same effects as the generalized active and inertia forces associated with blade 1. Thus,

$$F_r|_{Bl} = F_r^*|_{Bl} + F_r|_{AeroBl} + F_r|_{GravBl} + F_r|_{ElasticBl} + F_r|_{DampBl} \quad (r = 1, 2, \dots, 22)$$

Since ${}^E \mathbf{v}_r^C$ and ${}^E \boldsymbol{\omega}_r^H$ (and ${}^E \mathbf{v}_r^Q$ and ${}^E \mathbf{v}_r^{SI}(0)$) are equal to zero unless $r = 1, 2, \dots, 14; Teet$, the generalized active forces associated blade elasticity and damping do not contribute to the root loads (since also, $F_r|_{ElasticBl}$ and $F_r|_{DampBl}$ are equal to zero within this range of r 's). So,

$$F_r|_{Bl} = F_r^*|_{Bl} + F_r|_{AeroBl} + F_r|_{GravBl} \quad (r = 1, 2, \dots, 14; Teet)$$

Thus,

$$F_r|_{Bl} = \int_0^{BldFlexL} {}^E \mathbf{v}_r^{SI}(r) \cdot [\mathbf{F}_{AeroBl}^{SI}(r) - \mu^{Bl}(r) g \mathbf{z}_2 - \mu^{Bl}(r) {}^E \mathbf{a}^{SI}(r)] dr + \int_0^{BldFlexL} {}^E \boldsymbol{\omega}_r^{Ml}(r) \cdot \mathbf{M}_{AeroBl}^{Ml}(r) dr + {}^E \mathbf{v}_r^{SI}(BldFlexL) \cdot \{ \mathbf{F}_{TipDragBl}^{SI}(BldFlexL) - m^{BlTip} [g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL)] \} \quad (r = 1, 2, \dots, 14; Teet)$$

Now noting that ${}^E \mathbf{v}_r^{SI}(r) = {}^E \mathbf{v}_r^Q + {}^H \mathbf{v}_r^{SI}(r) + {}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QSI}(r)$, this can be expanded as follows:

$$F_r|_{Bl} = \int_0^{BldFlexL} [{}^E \mathbf{v}_r^Q + {}^H \mathbf{v}_r^{SI}(r)] \cdot [\mathbf{F}_{AeroBl}^{SI}(r) - \mu^{Bl}(r) g \mathbf{z}_2 - \mu^{Bl}(r) {}^E \mathbf{a}^{SI}(r)] dr + [{}^E \mathbf{v}_r^Q + {}^H \mathbf{v}_r^{SI}(BldFlexL)] \cdot \{ \mathbf{F}_{TipDragBl}^{SI}(BldFlexL) - m^{BlTip} [g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL)] \} + \int_0^{BldFlexL} [{}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QSI}(r)] \cdot [\mathbf{F}_{AeroBl}^{SI}(r) - \mu^{Bl}(r) g \mathbf{z}_2 - \mu^{Bl}(r) {}^E \mathbf{a}^{SI}(r)] dr + \int_0^{BldFlexL} {}^E \boldsymbol{\omega}_r^{Ml}(r) \cdot \mathbf{M}_{AeroBl}^{Ml}(r) dr + [{}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QSI}(BldFlexL)] \cdot \{ \mathbf{F}_{TipDragBl}^{SI}(BldFlexL) - m^{BlTip} [g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL)] \} \quad (r = 1, 2, \dots, 14; Teet)$$

However, since r is constrained to be between 1, 2, ..., 14; Teet and since ${}^H \mathbf{v}_r^{SI}(r)$ is equal to zero and ${}^E \boldsymbol{\omega}_r^{Ml}(r)$ equals ${}^E \boldsymbol{\omega}_r^H$ with this constraint, this can be simplified as follows:

$$\begin{aligned}
F_r|_{B1} = & \int_0^{BldFlexL} {}^E \mathbf{v}_r^Q \cdot \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr + {}^E \mathbf{v}_r^Q \cdot \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\} \\
& + \int_0^{BldFlexL} \left[{}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QS1}(r) \right] \cdot \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr + \int_0^{BldFlexL} {}^E \boldsymbol{\omega}_r^H \cdot \mathbf{M}_{AeroB1}^{M1}(r) dr \\
& + \left[{}^E \boldsymbol{\omega}_r^H \times \mathbf{r}^{QS1}(BldFlexL) \right] \cdot \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\}
\end{aligned} \quad (r = 1, 2, \dots, 14; Teet)$$

Or by engaging the cyclic permutation law of the scalar triple product,

$$\begin{aligned}
F_r|_{B1} = & \int_0^{BldFlexL} {}^E \mathbf{v}_r^Q \cdot \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr + {}^E \mathbf{v}_r^Q \cdot \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\} \\
& + \int_0^{BldFlexL} {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{r}^{QS1}(r) \times \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] \right\} dr + \int_0^{BldFlexL} {}^E \boldsymbol{\omega}_r^H \cdot \mathbf{M}_{AeroB1}^{M1}(r) dr \\
& + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{r}^{QS1}(BldFlexL) \times \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\} \right\}
\end{aligned} \quad (r = 1, 2, \dots, 14; Teet)$$

Thus it is seen that,

$$\mathbf{F}_{B1}^{SI}(0) = \int_0^{BldFlexL} \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr + \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right]$$

and

$$\begin{aligned}
\mathbf{M}_{B1}^H(0) + \mathbf{r}^{QS1}(0) \times \mathbf{F}_{B1}^{SI}(0) = & \int_0^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr + \int_0^{BldFlexL} \mathbf{r}^{QS1}(r) \times \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr \\
& + \mathbf{r}^{QS1}(BldFlexL) \times \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\}
\end{aligned}$$

or

$$\begin{aligned}
\mathbf{M}_{B1}^H(0) = & \int_0^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr + \int_0^{BldFlexL} \mathbf{r}^{QS1}(r) \times \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr \\
& + \mathbf{r}^{QS1}(BldFlexL) \times \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\} \\
& - \mathbf{r}^{QS1}(0) \times \left\{ \int_0^{BldFlexL} \left[\mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r) {}^E \mathbf{a}^{SI}(r) \right] dr + \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{SI}(BldFlexL) \right] \right\}
\end{aligned}$$

or

$$\mathbf{M}_{B1}^H(0) = \int_0^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr + \int_0^{BldFlexL} \left[\mathbf{r}^{QS1}(r) - \mathbf{r}^{QS1}(0) \right] \times \left[\mathbf{F}_{AeroB1}^{S1}(r) - \mu^{B1}(r) g\mathbf{z}_2 - \mu^{B1}(r) {}^E\mathbf{a}^{S1}(r) \right] dr \\ + \left[\mathbf{r}^{QS1}(BldFlexL) - \mathbf{r}^{QS1}(0) \right] \times \left\{ \mathbf{F}_{TipDragB1}^{S1}(BldFlexL) - m^{B1Tip} \left[g\mathbf{z}_2 + {}^E\mathbf{a}^{S1}(BldFlexL) \right] \right\}$$

Thus,

$$\mathbf{F}_{B1}^{S1}(0) = \int_0^{BldFlexL} \left\{ \mathbf{F}_{AeroB1}^{S1}(r) - \mu^{B1}(r) \left\{ g\mathbf{z}_2 + {}^E\mathbf{v}_I^{S1}(r) \ddot{q}_I + \left[\sum_{i=1}^{14} {}^E\mathbf{v}_i^{S1}(r) \ddot{q}_i \right] + \left[\sum_{i=16}^{18} {}^E\mathbf{v}_i^{S1}(r) \ddot{q}_i \right] + {}^E\mathbf{v}_{Teet}^{S1}(r) \ddot{q}_{Teet} \right\} \right. \\ \left. + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E\mathbf{v}_i^{S1}(r)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E\mathbf{v}_i^{S1}(r)) \dot{q}_i \right] + \frac{d}{dt} ({}^E\mathbf{v}_{Teet}^{S1}(r)) \dot{q}_{Teet} \right\} dr \\ + \mathbf{F}_{TipDragB1}^{S1}(BldFlexL) - m^{B1Tip} \left\{ g\mathbf{z}_2 + \left[\sum_{i=1}^{14} {}^E\mathbf{v}_i^{S1}(BldFlexL) \ddot{q}_i \right] + \left[\sum_{i=16}^{18} {}^E\mathbf{v}_i^{S1}(BldFlexL) \ddot{q}_i \right] + {}^E\mathbf{v}_{Teet}^{S1}(BldFlexL) \ddot{q}_{Teet} \right. \\ \left. + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E\mathbf{v}_i^{S1}(BldFlexL)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E\mathbf{v}_i^{S1}(BldFlexL)) \dot{q}_i \right] + \frac{d}{dt} ({}^E\mathbf{v}_{Teet}^{S1}(BldFlexL)) \dot{q}_{Teet} \right\}$$

and

$$\mathbf{M}_{B1}^H(0) = \int_0^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr$$

$$\begin{aligned}
 & + \int_0^{BldFlexL} [\mathbf{r}^{QS1}(r) - \mathbf{r}^{QS1}(0)] \times \left\{ \mathbf{F}_{AeroB1}^{S1}(r) - \mu^{B1}(r) \left\{ \begin{aligned} & g\mathbf{z}_2 + \left[\sum_{i=1}^{14} {}^E \mathbf{v}_i^{S1}(r) \ddot{q}_i \right] + \left[\sum_{i=16}^{18} {}^E \mathbf{v}_i^{S1}(r) \ddot{q}_i \right] + {}^E \mathbf{v}_{Teet}^{S1}(r) \ddot{q}_{Teet} \\ & + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{S1}(r)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{S1}(r)) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{S1}(r)) \dot{q}_{Teet} \end{aligned} \right\} \right\} dr \\
 & + [\mathbf{r}^{QS1}(BldFlexL) - \mathbf{r}^{QS1}(0)] \times \left\{ \mathbf{F}_{TipDragB1}^{S1}(BldFlexL) - m^{B1Tip} \left\{ \begin{aligned} & g\mathbf{z}_2 + \left[\sum_{i=1}^{14} {}^E \mathbf{v}_i^{S1}(BldFlexL) \ddot{q}_i \right] + \left[\sum_{i=16}^{18} {}^E \mathbf{v}_i^{S1}(BldFlexL) \ddot{q}_i \right] \\ & + {}^E \mathbf{v}_{Teet}^{S1}(BldFlexL) \ddot{q}_{Teet} \\ & + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{S1}(BldFlexL)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{S1}(BldFlexL)) \dot{q}_i \right] \\ & + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{S1}(BldFlexL)) \dot{q}_{Teet} \end{aligned} \right\} \right\}
 \end{aligned}$$

Or,

$$\begin{aligned}
 \mathbf{F}_{Bl_r}^{SI}(0) &= - \int_0^{BldFlexL} \mu^{B1}(r) {}^E \mathbf{v}_r^{SI}(r) dr - m^{B1Tip} {}^E \mathbf{v}_r^{SI}(BldFlexL) \quad (r = 1, 2, \dots, 14; 16, 17, 18; Teet) \\
 \mathbf{F}_{Bl_r}^{SI}(0) &= \int_0^{BldFlexL} \left\{ \mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) \left\{ g\mathbf{z}_2 + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(r)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(r)) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{SI}(r)) \dot{q}_{Teet} \right\} \right\} dr \\
 &\quad - m^{B1Tip} \left\{ g\mathbf{z}_2 + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(BldFlexL)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(BldFlexL)) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{SI}(BldFlexL)) \dot{q}_{Teet} \right\} \\
 &\quad + \mathbf{F}_{TipDragB1}^{SI}(BldFlexL)
 \end{aligned}$$

and

$$\begin{aligned}
 \mathbf{M}_{Bl_r}^H(0) &= - \int_0^{BldFlexL} [\mathbf{r}^{QS1}(r) - \mathbf{r}^{QS1}(0)] \times [\mu^{B1}(r) {}^E \mathbf{v}_r^{SI}(r)] dr - m^{B1Tip} [\mathbf{r}^{QS1}(BldFlexL) - \mathbf{r}^{QS1}(0)] \times {}^E \mathbf{v}_r^{SI}(BldFlexL) \quad (r = 1, 2, \dots, 14; 16, 17, 18; Teet) \\
 \mathbf{M}_{Bl_r}^H(0) &= \int_0^{BldFlexL} [\mathbf{r}^{QS1}(r) - \mathbf{r}^{QS1}(0)] \times \left\{ \mathbf{F}_{AeroB1}^{SI}(r) - \mu^{B1}(r) \left\{ g\mathbf{z}_2 + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(r)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(r)) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{SI}(r)) \dot{q}_{Teet} \right\} \right\} dr \\
 &\quad + [\mathbf{r}^{QS1}(BldFlexL) - \mathbf{r}^{QS1}(0)] \times \left\{ \mathbf{F}_{TipDragB1}^{SI}(BldFlexL) - m^{B1Tip} \left\{ g\mathbf{z}_2 + \left[\sum_{i=4}^{14} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(BldFlexL)) \dot{q}_i \right] + \left[\sum_{i=16}^{18} \frac{d}{dt} ({}^E \mathbf{v}_i^{SI}(BldFlexL)) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{Teet}^{SI}(BldFlexL)) \dot{q}_{Teet} \right\} \right\} \\
 &\quad + \int_0^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr
 \end{aligned}$$

The output loads are as follows:

$$RootFxc1 = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{i}_1^{B1} / 1,000$$

Blade 1 OoP shear force at the blade root (directed along the xc1-axis), (kN)

$$RootFyc1 = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{i}_2^{B1} / 1,000$$

Blade 1 IP shear force at the blade root (directed along the yc1-axis), (kN)

$$RootFxb1 = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{j}_1^{B1} / 1,000$$

Blade 1 flapwise shear force at the blade root (directed along the xb1-axis), (kN)

$$RootFyb1 = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{j}_2^{B1} / 1,000$$

Blade 1 edgewise shear force at the blade root (directed along the yb1-axis), (kN)

$$RootFzcl = RootFzbl = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{i}_3^{B1} / 1,000 = \mathbf{F}_{B1}^{SI}(0) \cdot \mathbf{j}_3^{B1} / 1,000$$

Blade 1 axial force at the blade root (directed along the zc1-/zb1-axis), (kN)

$$RootMxc1 = RootMIP1 = \mathbf{M}_{B1}^H(0) \cdot \mathbf{i}_1^{B1} / 1,000$$

Blade 1 IP moment (i.e., the moment caused by IP forces) at the blade root (about the xc1-axis),

(kN·m)

$$RootMyc1 = RootMOoP1 = \mathbf{M}_{B1}^H(0) \cdot \mathbf{i}_2^{B1} / 1,000$$

Blade 1 OoP moment (i.e., the moment caused by OoP forces) at the blade root (about the yc1-

axis), (kN·m)

$$RootMxb1 = RootMEdg1 = \mathbf{M}_{B1}^H(0) \cdot \mathbf{j}_1^{B1} / 1,000$$

Blade 1 edgewise moment (i.e., the moment caused by edgewise forces) at the blade root (about

the xb1-axis), (kN·m)

$$RootMyb1 = RootMFlp1 = \mathbf{M}_{B1}^H(0) \cdot \mathbf{j}_2^{B1} / 1,000$$

Blade 1 flapwise moment (i.e., the moment caused by flapwise forces) at the blade root (about the

yb1-axis), (kN·m)

$$RootMzcl = RootMzbl = \mathbf{M}_{B1}^H(0) \cdot \mathbf{i}_3^{B1} / 1,000 = \mathbf{M}_{B1}^H(0) \cdot \mathbf{j}_3^{B1} / 1,000$$

Blade 1 pitching moment at the blade root (about the zc1-/zb1-axis),

(kN·m)

Blade 1 Local Moment Outputs:

There are 3 output loads at any of the selected span stations i ($r = R^{Span\ i}$) of blade 1 ($i=1,2,\dots,5$). These are the 3 components of the bending moment $\mathbf{M}_{B1}^{M1}(R^{Span\ i})$ expressed in the *local* blade coordinate system (principal structural axes). Examining the results for the blade 1 root loads, it follows that:

$$\begin{aligned} \mathbf{M}_{B1}^{M1}(R^{Span\ i}) = & \int_{R^{Span\ i}}^{BldFlexL} \mathbf{M}_{AeroB1}^{M1}(r) dr + \int_{R^{Span\ i}}^{BldFlexL} \left[\mathbf{r}^{QS1}(r) - \mathbf{r}^{QS1}(R^{Span\ i}) \right] \times \left[\mathbf{F}_{AeroB1}^{S1}(r) - \mu^{B1}(r) g \mathbf{z}_2 - \mu^{B1}(r)^E \mathbf{a}^{S1}(r) \right] dr \\ & + \left[\mathbf{r}^{QS1}(BldFlexL) - \mathbf{r}^{QS1}(R^{Span\ i}) \right] \times \left\{ \mathbf{F}_{TipDragB1}^{S1}(BldFlexL) - m^{B1Tip} \left[g \mathbf{z}_2 + {}^E \mathbf{a}^{S1}(BldFlexL) \right] \right\} \quad (i = 1, 2, \dots, 5) \end{aligned}$$

The output loads are as follows:

$$SpniMLxb1 = \mathbf{M}_{B1}^{M1}(R^{Span\ i}) \cdot \mathbf{n}_1^{B1}(R^{Span\ i}) / 1,000$$

$$SpniMLyb1 = \mathbf{M}_{B1}^{M1}(R^{Span\ i}) \cdot \mathbf{n}_2^{B1}(R^{Span\ i}) / 1,000$$

$$SpniMLzb1 = \mathbf{M}_{B1}^{M1}(R^{Span\ i}) \cdot \mathbf{n}_3^{B1}(R^{Span\ i}) / 1,000$$

Blade 1 *local* edgewise moment at span station i (about the *local* xb1-structural axis), (kN·m)

Blade 1 *local* flapwise moment at span station i (about the *local* yb1-structural axis), (kN·m)

Blade 1 pitching moment at span station i (about the zc1-/zb1-/local zb1-axis), (kN·m)

Blade 2 Root Loads:

The equations for $\mathbf{F}_{B2}^{S2}(0)$, $\mathbf{M}_{B2}^H(0)$, $\mathbf{F}_{B2_r}^{S2}(0)$, $\mathbf{F}_{B2_t}^{S2}(0)$, $\mathbf{M}_{B2_r}^H(0)$, $\mathbf{M}_{B2_t}^H(0)$, and all 10 output loads are similar to blade 1.

Hub and Rotor Loads:

There are 14 output loads at the hub end of the low-speed shaft. 5 of them are the 3 components of the thrust and shear force F_{Rotor}^P (2 components are expressed in a nonrotating frame, 2 components are expressed in a rotating frame, and 1 component is independent of rotation). 5 other loads are the 3 components of the shaft bending moments, $M_{Rotor}^{L@P}$ (again, 2 components are expressed in a nonrotating frame, 2 components are expressed in a rotating frame, and 1 component is independent of rotation). The 11th and 12th loads are the rotor power and rotor power coefficient, respectively. The 13th and 14th loads are the rotor thrust and rotor torque coefficients, respectively. For a 2-blader, all these loads are given relative to the teeter pin (point P) as indicated. For the 3-blader, all of these loads are given relative to the apex of rotation (point Q, which is coincident with point P). The new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Rotor} = {}^E v_r^P \cdot F_{Rotor}^P + {}^E \omega_r^L \cdot M_{Rotor}^{L@P} \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with blade 1, blade 2, the hub, and the teeter springs and dampers. Thus,

$$\begin{aligned} F_r|_{Rotor} = & F_r^*|_{B1} + F_r|_{AeroB1} + F_r|_{GravB1} + F_r|_{ElasticB1} + F_r|_{DampB1} \\ & + F_r^*|_{B2} + F_r|_{AeroB2} + F_r|_{GravB2} + F_r|_{ElasticB2} + F_r|_{DampB2} \\ & + F_r^*|_H + F_r|_{GravH} + F_r|_{SpringTeet} + F_r|_{DampTeet} \end{aligned} \quad (r = 1, 2, \dots, 22)$$

Since ${}^E v_r^P$ and ${}^E \omega_r^L$ are equal to zero unless $r = 1, 2, \dots, 14$, the generalized active forces associated with blade and teeter elasticity and damping do not contribute to the hub and rotor loads (since also, $F_r|_{ElasticB1}$, $F_r|_{DampB1}$, $F_r|_{ElasticB2}$, $F_r|_{DampB2}$, $F_r|_{SpringTeet}$, and $F_r|_{DampTeet}$ are equal to zero if $r = 1, 2, \dots, 14$). So,

$$F_r|_{Rotor} = F_r^*|_{B1} + F_r|_{AeroB1} + F_r|_{GravB1} + F_r^*|_{B2} + F_r|_{AeroB2} + F_r|_{GravB2} + F_r^*|_H + F_r|_{GravH} \quad (r = 1, 2, \dots, 14)$$

When using the results for the blade 1 and blade 2 root loads, this equation can be simplified as follows:

$$F_r|_{Rotor} = F_r|_{B1} + F_r|_{B2} + F_r^*|_H + F_r|_{GravH} \quad (r = 1, 2, \dots, 14)$$

Thus,

$$F_r|_{Rotor} = {}^E \mathbf{v}_r^C \cdot \mathbf{F}_{B1}^{S1}(0) + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + [\mathbf{r}^{QS1}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B1}^{S1}(0) \right\} + {}^E \mathbf{v}_r^C \cdot \mathbf{F}_{B2}^{S2}(0) + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{QS2}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B2}^{S2}(0) \right\} \\ - m^H {}^E \mathbf{v}_r^C \cdot ({}^E \mathbf{a}^C + g\mathbf{z}_2) - {}^E \boldsymbol{\omega}_r^H \cdot (\bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H + {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H) \quad (r = 1, 2, \dots, 14)$$

or when grouping like terms:

$$F_r|_{Rotor} = {}^E \mathbf{v}_r^C \cdot \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H ({}^E \mathbf{a}^C + g\mathbf{z}_2) \right] \\ + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{QS1}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B1}^{S1}(0) + [\mathbf{r}^{QS2}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B2}^{S2}(0) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} \quad (r = 1, 2, \dots, 14)$$

Recognizing that

$${}^E \mathbf{v}_r^C = {}^E \mathbf{v}_r^P + {}^E \boldsymbol{\omega}_r^H \times (\mathbf{r}^{PQ} + \mathbf{r}^{QC})$$

, this generalized force can be expanded to:

$$F_r|_{Rotor} = \left[{}^E \mathbf{v}_r^P + {}^E \boldsymbol{\omega}_r^H \times (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \right] \cdot \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H ({}^E \mathbf{a}^C + g\mathbf{z}_2) \right] \\ + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{QS1}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B1}^{S1}(0) + [\mathbf{r}^{QS2}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B2}^{S2}(0) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} \quad (r = 1, 2, \dots, 14)$$

Now applying the cyclic permutation law of the scalar triple product:

$$F_r|_{Rotor} = {}^E \mathbf{v}_r^P \cdot \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H ({}^E \mathbf{a}^C + g\mathbf{z}_2) \right] + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \times \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H ({}^E \mathbf{a}^C + g\mathbf{z}_2) \right] \right\} \\ + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{QS1}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B1}^{S1}(0) + [\mathbf{r}^{QS2}(0) - \mathbf{r}^{QC}] \times \mathbf{F}_{B2}^{S2}(0) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} \quad (r = 1, 2, \dots, 14)$$

which simplifies to:

$$F_r|_{Rotor} = {}^E \mathbf{v}_r^P \cdot \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H ({}^E \mathbf{a}^C + g\mathbf{z}_2) \right] \\ + {}^E \boldsymbol{\omega}_r^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0)] \times \mathbf{F}_{B1}^{S1}(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0)] \times \mathbf{F}_{B2}^{S2}(0) \right. \\ \left. - m^H (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \times ({}^E \mathbf{a}^C + g\mathbf{z}_2) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} \quad (r = 1, 2, \dots, 14)$$

However, ${}^E \boldsymbol{\omega}_r^H$ equals ${}^E \boldsymbol{\omega}_r^L$ when r is not equal to Teet. Thus the generalized active force associated with the rotor can be expressed as follows:

$$F_r|_{Rotor} = {}^E \mathbf{v}_r^P \cdot \left[\mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H \left({}^E \mathbf{a}^C + g\mathbf{z}_2 \right) \right] \\ + {}^E \boldsymbol{\omega}_r^L \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0) \right] \times \mathbf{F}_{B1}^{S1}(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0) \right] \times \mathbf{F}_{B2}^{S2}(0) \right. \\ \left. - m^H \left(\mathbf{r}^{PQ} + \mathbf{r}^{QC} \right) \times \left({}^E \mathbf{a}^C + g\mathbf{z}_2 \right) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} \quad (r = 1, 2, \dots, 14)$$

Thus it is seen that,

$$\mathbf{F}_{Rotor}^P = \mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H \left({}^E \mathbf{a}^C + g\mathbf{z}_2 \right)$$

and

$$\mathbf{M}_{Rotor}^{L@P} = \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0) \right] \times \mathbf{F}_{B1}^{S1}(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0) \right] \times \mathbf{F}_{B2}^{S2}(0) - m^H \left(\mathbf{r}^{PQ} + \mathbf{r}^{QC} \right) \times \left({}^E \mathbf{a}^C + g\mathbf{z}_2 \right) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H$$

Thus,

$$\mathbf{F}_{Rotor}^P = \mathbf{F}_{B1}^{S1}(0) + \mathbf{F}_{B2}^{S2}(0) - m^H \left\{ \left(\sum_{i=1}^{14} {}^E \mathbf{v}_i^C \ddot{q}_i \right) + {}^E \mathbf{v}_{Teet}^C \ddot{q}_{Teet} + \left[\sum_{i=4}^{14} \frac{d}{dt} \left({}^E \mathbf{v}_i^C \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{Teet}^C \right) \dot{q}_{Teet} + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Rotor}^{L@P} = \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0) \right] \times \mathbf{F}_{B1}^{S1}(0) + \left[\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0) \right] \times \mathbf{F}_{B2}^{S2}(0) \\ - m^H \left(\mathbf{r}^{PQ} + \mathbf{r}^{QC} \right) \times \left\{ \left(\sum_{i=1}^{14} {}^E \mathbf{v}_i^C \ddot{q}_i \right) + {}^E \mathbf{v}_{Teet}^C \ddot{q}_{Teet} + \left[\sum_{i=4}^{14} \frac{d}{dt} \left({}^E \mathbf{v}_i^C \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{Teet}^C \right) \dot{q}_{Teet} + g\mathbf{z}_2 \right\} \\ - \bar{\mathbf{I}}^H \cdot \left\{ \left(\sum_{i=4}^{14} {}^E \boldsymbol{\omega}_i^H \ddot{q}_i \right) + {}^E \boldsymbol{\omega}_{Teet}^H \ddot{q}_{Teet} + \left[\sum_{i=7}^{14} \frac{d}{dt} \left({}^E \boldsymbol{\omega}_i^H \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \boldsymbol{\omega}_{Teet}^H \right) \dot{q}_{Teet} \right\} - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H$$

Or,

$$\mathbf{F}_{Rotor_r}^P = \mathbf{F}_{Bl_r}^{S1}(0) + \mathbf{F}_{B2_r}^{S2}(0) - m^H \mathbf{v}_r^C \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22)$$

$$\mathbf{F}_{Rotor_i}^P = \mathbf{F}_{Bl_i}^{S1}(0) + \mathbf{F}_{B2_i}^{S2}(0) - m^H \left\{ \left[\sum_{i=4}^{14} \frac{d}{dt} (\mathbf{v}_i^C) \dot{q}_i \right] + \frac{d}{dt} (\mathbf{v}_{Teet}^C) \dot{q}_{Teet} + g \mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Rotor_r}^{L@P} = \mathbf{M}_{Bl_r}^H(0) + \mathbf{M}_{B2_r}^H(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0)] \times \mathbf{F}_{Bl_r}^{S1}(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0)] \times \mathbf{F}_{B2_r}^{S2}(0) - m^H (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \times \mathbf{v}_r^C - \bar{\mathbf{I}}^H \cdot \mathbf{e}^H \omega_r^H \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22)$$

$$\mathbf{M}_{Rotor_i}^{L@P} = \mathbf{M}_{Bl_i}^H(0) + \mathbf{M}_{B2_i}^H(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0)] \times \mathbf{F}_{Bl_i}^{S1}(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0)] \times \mathbf{F}_{B2_i}^{S2}(0) - m^H (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \times \left\{ \left[\sum_{i=4}^{14} \frac{d}{dt} (\mathbf{v}_i^C) \dot{q}_i \right] + \frac{d}{dt} (\mathbf{v}_{Teet}^C) \dot{q}_{Teet} + g \mathbf{z}_2 \right\}$$

$$- \bar{\mathbf{I}}^H \cdot \left\{ \left[\sum_{i=7}^{14} \frac{d}{dt} (\mathbf{e}^H \omega_i^H) \dot{q}_i \right] + \frac{d}{dt} (\mathbf{e}^H \omega_{Teet}^H) \dot{q}_{Teet} \right\} - \mathbf{e}^H \omega^H \times \bar{\mathbf{I}}^H \cdot \mathbf{e}^H \omega^H$$

The output loads are as follows,

$$RotThrust = LSShftFxs = LSShftFxa = \mathbf{F}_{Rotor}^P \cdot \mathbf{e}_1 / 1,000 = \mathbf{F}_{Rotor}^P \cdot \mathbf{c}_1 / 1,000$$

Low-speed shaft thrust force (directed along the xs-/xa-axis) (this

is constant along the shaft and is equivalent to the rotor thrust force), (kN)

$$LSShftFya = \mathbf{F}_{Rotor}^P \cdot \mathbf{e}_2 / 1,000$$

Rotating low-speed shaft shear force (directed along the ya-axis) (this is constant along the shaft), (kN)

$$LSShftFza = \mathbf{F}_{Rotor}^P \cdot \mathbf{e}_3 / 1,000$$

Rotating low-speed shaft shear force (directed along the za-axis) (this is constant along the shaft), (kN)

$$LSShftFys = -\mathbf{F}_{Rotor}^P \cdot \mathbf{c}_3 / 1,000$$

Nonrotating low-speed shaft shear force (directed along the ys-axis) (this is constant along the shaft), (kN)

$$LSShftFzs = \mathbf{F}_{Rotor}^P \cdot \mathbf{c}_2 / 1,000$$

Nonrotating low-speed shaft shear force (directed along the zs-axis) (this is constant along the shaft), (kN)

$$RotTorq = LSShftTq = LSShftMxs = LSShftMxa = \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{e}_1 / 1,000 = \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{c}_1 / 1,000$$

Low-speed shaft torque (about the xs-/xa-axis) (this is

constant along the shaft and is equivalent to the rotor torque), (kN·m)

$$LSShftMya = \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{e}_2 / 1,000$$

Rotating low-speed shaft bending moment at the shaft tip [teeter pin for 2 blades] [apex of rotation for 3 blades]

(about the ya-axis), (kN·m)

$$LSShftMza = \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{e}_3 / 1,000$$

Rotating low-speed shaft bending moment at the shaft tip [teeter pin for 2 blades] [apex of rotation for 3 blades]

(about the za-axis), (kN·m)

$$LSShftMys = -\mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{c}_3 / 1,000$$

Nonrotating low-speed shaft bending moment at the shaft tip [teeter pin for 2 blades] [apex of rotation for 3

blades] (about the ys-axis), (kN·m)

$LSSTipMzs = \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{c}_2 / 1,000$ Nonrotating low-speed shaft bending moment at the shaft tip [teeter pin for 2 blades] [apex of rotation for 3 blades] (about the zs-axis), (kN·m)

$CThrstAzm = MOD \left[ATAN2(-CThrstzs, -CThrstys) \cdot \left(\frac{180}{\pi} \right) + 360 + AzimBIUp + 90, 360 \right]$ Azimuth location of the center of thrust (about the xs-/xa-axis), (deg)

$CThrstRad = CThrstArm = \frac{\sqrt{CThrstys^2 + CThrstzs^2}}{AvgNrmTpRd}$ Dimensionless radial (arm) location of the center of thrust (always positive, directly radially outboard at azimuth angle CThrstAzm), (-)

where: $CThrstys = -\frac{LSSTipMzs}{RotThrust}$ and $CThrstzs = \frac{LSSTipMys}{RotThrust}$

$RotPwr = LSShftPwr = (\dot{q}_{DrTr} + \dot{q}_{GeAz}) \cdot RotTorq = (\dot{q}_{DrTr} + \dot{q}_{GeAz}) \cdot LSShftTq$ Low-speed shaft power (this is equivalent to the rotor power), (kW)

$RotCp = LSShftCp = \frac{1,000 \cdot RotPwr}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^3}$ Low-speed shaft power coefficient (this is equivalent to the rotor power coefficient), (-)

$RotCq = LSShftCq = \frac{1,000 \cdot RotTorq}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^2 \cdot TipRad}$ Rotor torque coefficient, (-)

$RotCt = \frac{1,000 \cdot RotThrust}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^2}$ Rotor thrust coefficient, (-)

where V_0 is the hub-height wind speed and the projected area of the rotor, $ProjArea$, is found as follows:

$$ProjArea = \pi TipRad^2 \left\{ \frac{\cos[PreCone(1)] + \cos[PreCone(2)]}{2} \right\}^2$$

The rotor torque is equal to low-speed shaft torque as seen above. It is noted that this torque can be computed differently using the drivetrain flexibility and damping, though the load summation method and this other constraint method are equivalent. This can be demonstrated as follows. First of all, the equation above is equivalent to saying:

$$LSShftTq = {}^E \omega_{DrTr}^L \cdot \mathbf{M}_{Rotor}^{L@P} / 1,000$$

However, since ${}^E \mathbf{v}_{DrTr}^P$ is equal to zero, it is also equivalent to say:

$$LSShftTq = \left({}^E \mathbf{v}_{DrTr}^P \cdot \mathbf{F}_{Rotor}^P + {}^E \boldsymbol{\omega}_{DrTr}^L \cdot \mathbf{M}_{Rotor}^{L@P} \right) / 1,000$$

or,

$$LSShftTq = F_{DrTr|Rotor} / 1,000 \quad \text{or} \quad LSShftTq = \left(F_{DrTr|B1}^* + F_{DrTr|AeroB1} + F_{DrTr|GravB1} + F_{DrTr|B2}^* + F_{DrTr|AeroB2} + F_{DrTr|GravB2} + F_{DrTr|H}^* + F_{DrTr|GravH} \right) / 1,000$$

From the equations of motion, it is easily seen that this is equivalent to saying:

$$LSShftTq = \left(-F_{DrTr|ElasticDrive} - F_{DrTr|DampDrive} \right) / 1,000$$

and thus,

$$LSShftTq = \left(DTTorSpr \cdot q_{DrTr} + DTTorDmp \cdot \dot{q}_{DrTr} \right) / 1,000 \quad (= \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{c}_I / 1,000 \text{ and is equivalent to the rotor torque})$$

Thus, both the load summation method and the constraint method are equivalent. However, if the drivetrain DOF is disabled, then q_{DrTr} will equal zero and $\dot{q}_{DrTr}$ will equal zero, which implies that, at least, $DTTorSpr$ is equal to infinity (since the product of $DTTorSpr$ and q_{DrTr} is, in general, nonzero). Thus, to avoid using 2 different methods to calculate $LSShftTq$, it is best just to use $\mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{c}_I / 1,000$, which will always work, regardless of the number of DOFs disabled.

Like the $LSShftTq$, it is noted that $LSSTipMya$ can also be computed differently using the teeter springs and dampers, though the load summation method and this other constraint method are equivalent. This also can be demonstrated as follows. First of all, the equation above is equivalent to saying:

$$LSSTipMya = {}^E \boldsymbol{\omega}_{Teet}^H \cdot \mathbf{M}_{Rotor}^{L@P} / 1,000$$

Or,

$$LSSTipMya = {}^E \boldsymbol{\omega}_{Teet}^H \cdot \left\{ \mathbf{M}_{B1}^H(0) + \mathbf{M}_{B2}^H(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS1}(0)] \times \mathbf{F}_{B1}^{S1}(0) + [\mathbf{r}^{PQ} + \mathbf{r}^{QS2}(0)] \times \mathbf{F}_{B2}^{S2}(0) \right. \\ \left. - m^H (\mathbf{r}^{PQ} + \mathbf{r}^{QC}) \times ({}^E \mathbf{a}^C + g\mathbf{z}_2) - \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\alpha}^H - {}^E \boldsymbol{\omega}^H \times \bar{\mathbf{I}}^H \cdot {}^E \boldsymbol{\omega}^H \right\} / 1,000$$

Now applying the cyclic permutation law of the scalar triple product:

$$LSSTipMya = \left\{ \begin{aligned} & {}^E \omega_{Teet}^H \times [r^{PQ} + r^{QS1}(0)] \cdot F_{B1}^{S1}(0) + {}^E \omega_3^H \times [r^{PQ} + r^{QS2}(0)] \cdot F_{B2}^{S2}(0) - m^H {}^E \omega_{Teet}^H \times (r^{PQ} + r^{QC}) \cdot ({}^E a^C + gz_2) \\ & + {}^E \omega_{Teet}^H \cdot [M_{B1}^H(0) + M_{B2}^H(0) - \bar{I}^H \cdot {}^E a^H - {}^E \omega^H \times \bar{I}^H \cdot {}^E \omega^H] \end{aligned} \right\} / 1,000$$

Recognizing also that ${}^E v_{Teet}^{S1}(0) = {}^E \omega_{Teet}^H \times [r^{PQ} + r^{QS1}(0)]$, ${}^E v_{Teet}^{S2}(0) = {}^E \omega_{Teet}^H \times [r^{PQ} + r^{QS2}(0)]$, and ${}^E v_{Teet}^C = {}^E \omega_{Teet}^H \times (r^{PQ} + r^{QC})$, this can be simplified as follows:

$$LSSTipMya = \left[\begin{aligned} & {}^E v_{Teet}^{S1}(0) \cdot F_{B1}^{S1}(0) + {}^E \omega_{Teet}^H \cdot M_{B1}^H(0) + {}^E v_{Teet}^{S2}(0) \cdot F_{B2}^{S2}(0) + {}^E \omega_{Teet}^H \cdot M_{B2}^H(0) - m^H {}^E v_{Teet}^C \cdot ({}^E a^C + gz_2) \\ & - {}^E \omega_{Teet}^H \cdot (\bar{I}^H \cdot {}^E a^H + {}^E \omega^H \times \bar{I}^H \cdot {}^E \omega^H) \end{aligned} \right] / 1,000$$

or

$$LSSTipMya = (F_{Teet}|_{B1} + F_{Teet}|_{B2} + F_{Teet}^*|_H + F_{Teet}|_{GravH}) / 1,000$$

or,

$$LSSTipMya = (F_{Teet}^*|_H + F_{Teet}^*|_{B1} + F_{Teet}^*|_{B2} + F_{Teet}|_{AeroB1} + F_{Teet}|_{AeroB2} + F_{Teet}|_{GravH} + F_{Teet}|_{GravB1} + F_{Teet}|_{GravB2}) / 1,000$$

From the equations of motion, it is easily seen that this is equivalent to saying:

$$LSSTipMya = (-F_{Teet}|_{SpringTeet} - F_{Teet}|_{DampTeet}) / 1,000$$

and thus,

$$LSSTipMya = \left\{ \begin{aligned} & IF[|q_{Teet}| > TeetSSStP, TeetSSSp \cdot SIGN(q_{Teet})(|q_{Teet}| - TeetSSStP), 0] \\ & + IF[|q_{Teet}| > TeetHStP, TeetHSSp \cdot SIGN(q_{Teet})(|q_{Teet}| - TeetHStP), 0] \\ & + IF[\dot{q}_{Teet} < 0, TeetCDmp \cdot SIGN(\dot{q}_{Teet}), 0] + IF[|q_{Teet}| > TeetDmpP, TeetDmp \cdot \dot{q}_{Teet}, 0] \end{aligned} \right\} / 1,000 \quad (= \mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{e}_2 / 1,000)$$

Thus, both the load summation method and the constraint method are equivalent. Thus, to avoid using 2 different methods to calculate $LSSTipMya$ if various DOFs are disabled, it is best just to use $\mathbf{M}_{Rotor}^{L@P} \cdot \mathbf{e}_2 / 1,000$, which will always work.

Shaft Strain Gage Loads:

There are 4 output loads at point SG on the low-speed shaft [which is a point on the shaft a distance $ShftGagL$ towards the nacelle from point P (or point Q for a 3-blader since point P does not exist)]. These are 2 of the 3 components of the shaft bending moments, $M_{Rotor}^{L@SG}$ (2 components are expressed in a nonrotating frame, 2 components are expressed in a rotating frame, and third component which is directed in the c_1 direction is not used because it is the same as the rotor torque). Since the low-speed shaft is assumed to be rigid and massless between points P and SG, it is easily seen that:

$$M_{Rotor}^{L@SG} = M_{Rotor}^{L@P} - r^{PSG} \times F_{Rotor}^P$$

$$\text{since } r^{PSG} \text{ equals } -r^{SGP}.$$

Thus,

$$LSSGagMya = M_{Rotor}^{L@F} \cdot e_2 / 1,000 = LSSTipMya + ShftGagL \cdot LSShftFza$$

(about the ya-axis), (kN·m)

Rotating low-speed shaft bending moment at the shaft's strain gages

$$LSSGagMza = M_{Rotor}^{L@F} \cdot e_3 / 1,000 = LSSTipMza - ShftGagL \cdot LSShftFya$$

(about the za-axis), (kN·m)

Rotating low-speed shaft bending moment at the shaft's strain gages

$$LSSGagMys = -M_{Rotor}^{L@F} \cdot c_3 / 1,000 = LSSTipMys + ShftGagL \cdot LSShftFzs$$

(about the ys-axis), (kN·m)

Nonrotating low-speed shaft bending moment at the shaft's strain gages

$$LSSGagMzs = M_{Rotor}^{L@F} \cdot c_2 / 1,000 = LSSTipMzs - ShftGagL \cdot LSShftFys$$

(about the zs-axis), (kN·m)

Nonrotating low-speed shaft bending moment at the shaft's strain gages

Note that no shear or thrust forces need be output at point SG since these would be the same as the shear and thrust forces at point P. Note also that

$$c_1 \cdot M_{Rotor}^{L@P} = c_1 \cdot M_{Rotor}^{L@SG}$$

and thus the low-speed shaft torque or rotor torque are constant along the shaft.

Generator and High-Speed Shaft Loads:

There are 9 output loads on the high-speed shaft. The first and second are the high-speed shaft torque, $HSShftTq$, and high-speed shaft torque coefficient, $HSShftCq$, whose convention is that it has a positive value when the $LSShftTq$ is positive. The third and fourth are the high-speed shaft power, $HSShftPwr$, and high-speed shaft power coefficient, $HSShftCp$. The fifth and sixth are the generator electrical torque, $GenTq$, and generator electrical torque coefficient, $GenCq$. The seventh is the high-speed shaft braking torque, $HSSBrTq$. The eighth is the generator electrical power, $GenPwr$. The ninth is the electrical generator power coefficient, $GenCp$.

From a simple free-body diagram of a black-box gearbox,

$$HSShftTq = \frac{LSShftTq \cdot GBoxEff^{SIGN(LSShftTq)}}{GBRatio}$$

High-speed shaft torque (this is constant along the shaft and has the convention that it is positive

when the $LSShftTq$ is positive), (kN·m)

This can alternatively be written in terms of the high-speed shaft motions and torques through use of the equation for the GeAz DOF as follows.

From earlier work,

$$HSShftTq = \frac{{}^E\omega_{DrTr}^L \cdot M_{Rotor}^{L@P} GBoxEff^{SIGN(LSShftTq)}}{1,000 \cdot GBRatio}$$

or,

$$HSShftTq = \frac{\left(F_{DrTr}^*|_{B1} + F_{DrTr}|_{AeroB1} + F_{DrTr}|_{GravB1} + F_{DrTr}^*|_{B2} + F_{DrTr}|_{AeroB2} + F_{DrTr}|_{GravB2} + F_{DrTr}^*|_H + F_{DrTr}|_{GravH} \right) GBoxEff^{SIGN(LSShftTq)}}{1,000 \cdot GBRatio}$$

or,

$$HSShftTq = \frac{\left(F_{GeAz}^*|_{B1} + F_{GeAz}|_{AeroB1} + F_{GeAz}|_{GravB1} + F_{GeAz}^*|_{B2} + F_{GeAz}|_{AeroB2} + F_{GeAz}|_{GravB2} + F_{GeAz}^*|_H + F_{GeAz}|_{GravH} \right) GBoxEff^{SIGN(LSShftTq)}}{1,000 \cdot GBRatio}$$

From the equations of motion for the GeAz DOF, it is seen that this is equivalent to saying:

$$HSShftTq = \frac{\left(-F_{GeAz}^*|_G - F_{GeAz}|_{Gen} - F_{GeAz}|_{Brake} - F_{GeAz}|_{GBFric} \right) GBoxEff^{SIGN(LSShftTq)}}{1,000 \cdot GBRatio}$$

and thus,

$$HSShftTq = \frac{\left(\frac{GenIner \cdot GBRatio^2 \cdot \ddot{q}_{GeAz} + GenDir \cdot GenIner \cdot GBRatio \left\{ \left(\sum_{i=4}^{12} {}^E \omega_i^R \ddot{q}_i \right) + \left[\sum_{i=7}^{12} \frac{d}{dt} ({}^E \omega_i^R) \dot{q}_i \right] \right\} \cdot c_1}{GBoxEff^{SIGN(LSShftTq)}} + GBRatio \cdot T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t) + GBRatio \cdot T^{Brake}(t) \right) GBoxEff^{SIGN(LSShftTq)}}{1,000 \cdot GBRatio}$$

or,

$$HSShftTq = \frac{GenIner \cdot GBRatio \cdot \ddot{q}_{GeAz} + GenDir \cdot GenIner \cdot \left\{ \left(\sum_{i=4}^{12} {}^E \omega_i^R \ddot{q}_i \right) + \left[\sum_{i=7}^{12} \frac{d}{dt} ({}^E \omega_i^R) \dot{q}_i \right] \right\} \cdot c_1 + T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t) + T^{Brake}(t)}{1,000}$$

or,

$$HSShftTq = \left[GenIner \cdot GBRatio \cdot \ddot{q}_{GeAz} + GenDir \cdot GenIner \cdot {}^E \alpha^R \cdot c_1 + T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t) + T^{Brake}(t) \right] / 1,000$$

$$HSShftCq = \frac{1,000 \cdot HSShftTq}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^2 \cdot TipRad} \quad \text{High-speed shaft torque coefficient, (-)}$$

$$HSShftPwr = HSShftTq \cdot GBRatio \cdot \dot{q}_{GeAz} \quad \text{High-speed shaft power, (kW)}$$

$$HSShftCp = \frac{1,000 \cdot HSShftPwr}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^3} \quad \text{High-speed shaft power coefficient, (-)}$$

$$HSSBrTq = T^{Brake}(t) / 1,000 \quad \text{High-speed shaft braking torque, (kN·m)}$$

$$GenTq = T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t) / 1,000 \quad \text{Electrical generator torque (positive reflects power extracted and negative represents a motoring-up situation or power input), (kN·m)}$$

$$GenCq = \frac{1,000 \cdot GenTq}{\frac{1}{2} Rho \cdot ProjArea \cdot V_0^2 \cdot TipRad}$$

Electrical generator torque coefficient, (-)

Though the $HSShftTq$ is calculated the same regardless of the generator model employed, $GenPwr$ is not. Similar to how power is transmitted through the gearbox with a simple efficiency, for the simple generator or simple variable-speed generator control models, the electrical generator power is as follows:

$$GenPwr = GBRatio \cdot \dot{q}_{GeAz} \cdot GenTq \cdot GenEff^{SIGN[T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t)]} / 1,000$$

Electrical generator power (positive reflects power extracted and negative

represents a motoring-up situation or power input), (kW)

And for the Thevenin-Equivalent induction generator model,

$$GenPwr = (Pwr_{Mechanical} - Pwr_{StatorLoss} - Pwr_{ResistiveLoss}) / 1,000$$

Electrical generator power (positive reflects power extracted and negative

represents a motoring-up situation or power input), (kW)

where,

$$Pwr_{Mechanical} = GBRatio \cdot \dot{q}_{GeAz} \cdot T^{Gen}(GBRatio \cdot \dot{q}_{GeAz}, t)$$

(the sign of this is governed by T^{Gen})

$$Pwr_{StatorLoss} = TEC_NPha |\bar{I}_1|^2 TEC_SRes$$

(always positive)

and

$$Pwr_{ResistiveLoss} = TEC_NPha |\bar{I}_2|^2 TEC_RRes$$

(always positive)

where,

$$\bar{I}_2 = \frac{V_{IA}}{\left(R_{el} - \frac{TEC_RRes}{Slip} \right) + (X_{el} + TEC_RLR) \bar{j}}$$

and

$$\bar{I}_1 = \bar{I}_2 + \frac{V_{IA}}{TEC_MR\bar{j}}$$

where the definition of V_{IA} , R_{el} , X_{el} , and $Slip$ are given elsewhere and $\bar{j} = \sqrt{-1}$.

Otherwise, the electrical generator power, $GenPwr$, is a user-defined function of the high-speed shaft speed, $GBRatio$, $\dot{q}_{GeA}$, and time t .

Finally,

$$GenCp = \frac{1,000 \cdot GenPwr}{\frac{1}{2} \rho \cdot ProjArea \cdot V_0^3}$$

Electrical generator power coefficient, (-)

Rotor-Furl Axis Loads:

There is 1 output load on the rotor-furl axis. This is the rotor-furl moment about the rotor-furl axis. Of course, we could also output all 6 components of the force $\mathbf{F}_{Gen,Rot}^V$ / moment $\mathbf{M}_{Gen,Rot}^{N@V}$ acting on the rotor-furl axis at point V on the nacelle. Following the analysis for finding the blade root loads, the new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Gen,Rot} = {}^E \mathbf{v}_r^V \cdot \mathbf{F}_{Gen,Rot}^V + {}^E \boldsymbol{\omega}_r^N \cdot \mathbf{M}_{Gen,Rot}^{N@V} \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with blade 1, blade 2, the hub, the drivetrain, and the structure that furls with the rotor. Thus,

$$\begin{aligned} F_r|_{Gen,Rot} = & F_r|_{B1}^* + F_r|_{AeroB1} + F_r|_{GravB1} + F_r|_{ElasticB1} + F_r|_{DampB1} + F_r|_{B2}^* + F_r|_{AeroB2} + F_r|_{GravB2} + F_r|_{ElasticB2} + F_r|_{DampB2} \\ & + F_r|_H^* + F_r|_{GravH} + F_r|_{SpringTeet} + F_r|_{DampTeet} + F_r|_G^* + F_r|_{Gen} + F_r|_{Brake} + F_r|_{GBFric} + F_r|_{ElasticDrive} + F_r|_{DampDrive} \quad (r = 1, 2, \dots, 22) \\ & + F_r|_R^* + F_r|_{GravR} + F_r|_{SpringRF} + F_r|_{DampRF} \end{aligned}$$

Since ${}^E \mathbf{v}_r^V$ and ${}^E \boldsymbol{\omega}_r^N$ are equal to zero unless $r = 1, 2, \dots, 11$, the generalized active forces associated with blade, drivetrain, rotor-furl, and teeter elasticity and damping, as well as the generator torque, HSS braking torque, and gearbox friction do not contribute to the rotor-furl loads (since also, $F_r|_{ElasticB1}$, $F_r|_{DampB1}$, $F_r|_{ElasticB2}$, $F_r|_{DampB2}$, $F_r|_{SpringTeet}$, $F_r|_{DampTeet}$, $F_r|_{SpringRF}$, $F_r|_{DampRF}$, $F_r|_{ElasticDrive}$, $F_r|_{DampDrive}$, $F_r|_{Gen}$, $F_r|_{Brake}$, and $F_r|_{GBFric}$ are equal to zero if $r = 1, 2, \dots, 11$). So,

$$F_r|_{Gen,Rot} = F_r|_{B1}^* + F_r|_{AeroB1} + F_r|_{GravB1} + F_r|_{B2}^* + F_r|_{AeroB2} + F_r|_{GravB2} + F_r|_H^* + F_r|_{GravH} + F_r|_R^* + F_r|_{GravR} + F_r|_G^* \quad (r = 1, 2, \dots, 11)$$

When using the results for hub and rotor loads, this equation can be simplified as follows:

$$F_r|_{Gen,Rot} = F_r|_{Rotor} + F_r|_R^* + F_r|_{GravR} + F_r|_G^* \quad (r = 1, 2, \dots, 11)$$

Thus,

$$F_r|_{Gen,Rot} = {}^E \mathbf{v}_r^P \cdot \mathbf{F}_{Rotor}^P + {}^E \boldsymbol{\omega}_r^L \cdot \mathbf{M}_{Rotor}^{L@P} - m^R {}^E \mathbf{v}_r^D \cdot ({}^E \mathbf{a}^D + g \mathbf{z}_2) - {}^E \boldsymbol{\omega}_r^R \cdot (\bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R + {}^E \boldsymbol{\omega}_r^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}_r^R) - {}^E \boldsymbol{\omega}_r^G \cdot (\bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G + {}^E \boldsymbol{\omega}_r^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}_r^G) \quad (r = 1, 2, \dots, 11)$$

However, ${}^E \boldsymbol{\omega}_r^L$, ${}^E \boldsymbol{\omega}_r^G$, ${}^E \boldsymbol{\omega}_r^R$, and ${}^E \boldsymbol{\omega}_r^N$ are all equal when r is constrained to be between 1 and 11. Thus, when grouping like terms:

$$F_r|_{Gen,Rot} = {}^E \mathbf{v}_r^P \cdot \mathbf{F}_{Rotor}^P - m^R {}^E \mathbf{v}_r^D \cdot ({}^E \mathbf{a}^D + g\mathbf{z}_2) + {}^E \boldsymbol{\omega}_r^N \cdot \left(\mathbf{M}_{Rotor}^{L@P} - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \right) \quad (r = 1, 2, \dots, 11)$$

Recognizing also that ${}^E \mathbf{v}_r^P = {}^E \mathbf{v}_r^V + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{VP}$ and ${}^E \mathbf{v}_r^D = {}^E \mathbf{v}_r^V + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{VD}$, when $r = 1, 2, \dots, 11$, this generalized force can be expanded to:

$$F_r|_{Gen,Rot} = ({}^E \mathbf{v}_r^V + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{VP}) \cdot \mathbf{F}_{Rotor}^P - m^R ({}^E \mathbf{v}_r^V + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{VD}) \cdot ({}^E \mathbf{a}^D + g\mathbf{z}_2) + {}^E \boldsymbol{\omega}_r^N \cdot \left(\mathbf{M}_{Rotor}^{L@P} - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \right) \quad (r = 1, 2, \dots, 11)$$

Now applying the cyclic permutation law of the scalar triple product:

$$F_r|_{Gen,Rot} = {}^E \mathbf{v}_r^V \cdot \left[\mathbf{F}_{Rotor}^P - m^R ({}^E \mathbf{a}^D + g\mathbf{z}_2) \right] + {}^E \boldsymbol{\omega}_r^N \cdot \left[\mathbf{r}^{VP} \times \mathbf{F}_{Rotor}^P - m^R \mathbf{r}^{VD} \times ({}^E \mathbf{a}^D + g\mathbf{z}_2) \right] + {}^E \boldsymbol{\omega}_r^N \cdot \left(\mathbf{M}_{Rotor}^{L@P} - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \right) \quad (r = 1, 2, \dots, 11)$$

which simplifies to:

$$F_r|_{Gen,Rot} = {}^E \mathbf{v}_r^V \cdot \left[\mathbf{F}_{Rotor}^P - m^R ({}^E \mathbf{a}^D + g\mathbf{z}_2) \right] + {}^E \boldsymbol{\omega}_r^N \cdot \left(\mathbf{M}_{Rotor}^{L@P} + \mathbf{r}^{VP} \times \mathbf{F}_{Rotor}^P - m^R \mathbf{r}^{VD} \times ({}^E \mathbf{a}^D + g\mathbf{z}_2) - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \right) \quad (r = 1, 2, \dots, 11)$$

Thus it is seen that:

$$\mathbf{F}_{Gen,Rot}^V = \mathbf{F}_{Rotor}^P - m^R ({}^E \mathbf{a}^D + g\mathbf{z}_2)$$

and

$$\mathbf{M}_{Gen,Rot}^{N@V} = \mathbf{M}_{Rotor}^{L@P} + \mathbf{r}^{VP} \times \mathbf{F}_{Rotor}^P - m^R \mathbf{r}^{VD} \times ({}^E \mathbf{a}^D + g\mathbf{z}_2) - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\alpha}^R - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\alpha}^G - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G$$

Thus,

$$\mathbf{F}_{Gen,Rot}^V = \mathbf{F}_{Rotor}^P - m^R \left\{ \left(\sum_{i=1}^{12} {}^E \mathbf{v}_i^D \ddot{q}_i \right) + \left[\sum_{i=4}^{12} \frac{d}{dt} ({}^E \mathbf{v}_i^D) \dot{q}_i \right] + g \mathbf{z}_2 \right\}$$

and

$$\begin{aligned} \mathbf{M}_{Gen,Rot}^{N@V} = & \mathbf{M}_{Rotor}^{L@P} + \mathbf{r}^{VP} \times \mathbf{F}_{Rotor}^P - m^R \mathbf{r}^{VD} \times \left\{ \left(\sum_{i=1}^{12} {}^E \mathbf{v}_i^D \ddot{q}_i \right) + \left[\sum_{i=4}^{12} \frac{d}{dt} ({}^E \mathbf{v}_i^D) \dot{q}_i \right] + g \mathbf{z}_2 \right\} \\ & - \bar{\mathbf{I}}^R \cdot \left\{ \left(\sum_{i=4}^{12} {}^E \boldsymbol{\omega}_i^R \ddot{q}_i \right) + \left[\sum_{i=7}^{12} \frac{d}{dt} ({}^E \boldsymbol{\omega}_i^R) \dot{q}_i \right] \right\} - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot \left\{ \left(\sum_{i=4}^{13} {}^E \boldsymbol{\omega}_i^G \ddot{q}_i \right) + \left[\sum_{i=7}^{13} \frac{d}{dt} ({}^E \boldsymbol{\omega}_i^G) \dot{q}_i \right] \right\} - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \end{aligned}$$

Or,

$$\begin{aligned} \mathbf{F}_{Gen,Rot_r}^V &= \mathbf{F}_{Rotor_r}^P - m^R {}^E \mathbf{v}_r^D \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22) \\ \mathbf{F}_{Gen,Rot_r}^V &= \mathbf{F}_{Rotor_r}^P - m^R \left\{ \left[\sum_{i=4}^{12} \frac{d}{dt} ({}^E \mathbf{v}_i^D) \dot{q}_i \right] + g \mathbf{z}_2 \right\} \end{aligned}$$

and

$$\begin{aligned} \mathbf{M}_{Gen,Rot_r}^{N@V} &= \mathbf{M}_{Rotor_r}^{L@P} + \mathbf{r}^{VP} \times \mathbf{F}_{Rotor_r}^P - m^R \mathbf{r}^{VD} \times {}^E \mathbf{v}_r^D - \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}_r^R - \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}_r^G \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22) \\ \mathbf{M}_{Gen,Rot_r}^{N@V} &= \mathbf{M}_{Rotor_r}^{L@P} + \mathbf{r}^{VP} \times \mathbf{F}_{Rotor_r}^P - m^R \mathbf{r}^{VD} \times \left\{ \left[\sum_{i=4}^{12} \frac{d}{dt} ({}^E \mathbf{v}_i^D) \dot{q}_i \right] + g \mathbf{z}_2 \right\} \\ &\quad - \bar{\mathbf{I}}^R \cdot \left[\sum_{i=7}^{12} \frac{d}{dt} ({}^E \boldsymbol{\omega}_i^R) \dot{q}_i \right] - {}^E \boldsymbol{\omega}^R \times \bar{\mathbf{I}}^R \cdot {}^E \boldsymbol{\omega}^R - \bar{\mathbf{I}}^G \cdot \left[\sum_{i=7}^{13} \frac{d}{dt} ({}^E \boldsymbol{\omega}_i^G) \dot{q}_i \right] - {}^E \boldsymbol{\omega}^G \times \bar{\mathbf{I}}^G \cdot {}^E \boldsymbol{\omega}^G \end{aligned}$$

The output loads is as follows,

$$RFrIBrM = \mathbf{M}_{Gen,Rot}^{N@V} \cdot rfa / 1,000$$

Rotor-furl bearing furl moment (about the rotor-furl axis), (kN·m)

Like the *LSShftTq* and *LSSTipMza*, it is noted that the rotor-furling furl moment can be computed differently using the rotor-furl springs and dampers, though the load summation method and this other constraint method are equivalent. This can be demonstrated as follows. First of all, the equation above is equivalent to saying:

$$RFrIBrM = {}^E \boldsymbol{\omega}_{RFrl}^R \cdot \mathbf{M}_{Gen,Rot}^{N@V} / 1,000$$

Or,

$$RFrlBrM = {}^E \omega_{RFrl}^R \cdot \left[M_{Rotor}^{L@P} + r^{VP} \times F_{Rotor}^P - m^R r^{VD} \times ({}^E a^D + gz_2) - \bar{I}^R \cdot {}^E \alpha^R - {}^E \omega^R \times \bar{I}^R \cdot {}^E \omega^R - \bar{I}^G \cdot {}^E \alpha^G - {}^E \omega^G \times \bar{I}^G \cdot {}^E \omega^G \right] / 1,000$$

Now applying the cyclic permutation law of the scalar triple product:

$$RFrlBrM = \left[\begin{aligned} & {}^E \omega_{RFrl}^R \times r^{VP} \cdot F_{Rotor}^P - m^R {}^E \omega_{RFrl}^R \times r^{VD} \cdot ({}^E a^D + gz_2) \\ & + {}^E \omega_{RFrl}^R \cdot \left(M_{Rotor}^{L@P} - \bar{I}^R \cdot {}^E \alpha^R - {}^E \omega^R \times \bar{I}^R \cdot {}^E \omega^R - \bar{I}^G \cdot {}^E \alpha^G - {}^E \omega^G \times \bar{I}^G \cdot {}^E \omega^G \right) \end{aligned} \right] / 1,000$$

Recognizing also that ${}^E v_{RFrl}^P = {}^E \omega_{RFrl}^R \times r^{VP}$ and ${}^E v_{RFrl}^D = {}^E \omega_{RFrl}^R \times r^{VD}$, and also that ${}^E \omega_{RFrl}^L$, ${}^E \omega_{RFrl}^R$, and ${}^E \omega_{RFrl}^G$, are equal, this can be expanded as follows:

$$RFrlBrM = \left[{}^E v_{RFrl}^P \cdot F_{Rotor}^P + {}^E \omega_{RFrl}^L \cdot M_{Rotor}^{L@P} - m^R {}^E v_{RFrl}^D \cdot ({}^E a^D + gz_2) - {}^E \omega_{RFrl}^R \cdot \left(\bar{I}^R \cdot {}^E \alpha^R + {}^E \omega^R \times \bar{I}^R \cdot {}^E \omega^R \right) - {}^E \omega_{RFrl}^G \cdot \left(\bar{I}^G \cdot {}^E \alpha^G + {}^E \omega^G \times \bar{I}^G \cdot {}^E \omega^G \right) \right] / 1,000$$

or,

$$RFrlBrM = \left(F_{RFrl}|_{Rotor} + F_{RFrl}|_R^* + F_{RFrl}|_{GravR} + F_{RFrl}|_G^* \right) / 1,000$$

or,

$$RFrlBrM = \left(F_{RFrl}|_R^* + F_{RFrl}|_H^* + F_{RFrl}|_{B1}^* + F_{RFrl}|_{B2}^* + F_{RFrl}|_G^* + F_{RFrl}|_{AeroB1} + F_{RFrl}|_{AeroB2} + F_{RFrl}|_{GravR} + F_{RFrl}|_{GravH} + F_{RFrl}|_{GravB1} + F_{RFrl}|_{GravB2} \right) / 1,000$$

From the equations of motion, it is easily seen that this is equivalent to saying:

$$RFrlBrM = \left(-F_{RFrl}|_{SpringRF} - F_{RFrl}|_{DampRF} \right) / 1,000$$

and thus,

$$RFrIBrM = \left\{ \begin{array}{l} RFrISpr \cdot q_{RFl} + IF[q_{RFl} > RFrIUSSP, RFrIUSSpr(q_{RFl} - RFrIUSSP), 0] \\ + IF[q_{RFl} < RFrIDSSP, RFrIDSSpr(q_{RFl} - RFrIDSSP), 0] \\ + RFrIDmp \cdot \dot{q}_{RFl} + IF[\dot{q}_{RFl} < 0, RFrICDmp \cdot SIGN(\dot{q}_{RFl}), 0] \\ + IF[q_{RFl} > RFrIUSDP, RFrIUSDmp \cdot \dot{q}_{RFl}, 0] + IF[q_{RFl} < RFrIDSDP, RFrIDSDmp \cdot \dot{q}_{RFl}, 0] \end{array} \right\} / 1,000 \quad (= \mathbf{M}_{Gen, Rot}^{N@V} \cdot \mathbf{rfa} / 1,000)$$

Thus, both the load summation method and the constraint method are equivalent. Thus, to avoid using 2 different methods to calculate $RFrIBrM$ if various DOFs are disabled, it is best just to use $\mathbf{M}_{Gen, Rot}^{N@V} \cdot \mathbf{rfa} / 1,000$, which will always work.

Tail-Furl Axis Loads:

There is 1 output load on the tail-furl axis. This is the tail-furl moment about the tail-furl axis. Of course, we could also output all 6 components of the force $\mathbf{F}_{Tail}^W$ / moment $\mathbf{M}_{Tail}^{N@W}$ acting on the tail-furl axis at point W on the nacelle. Following the analysis for finding the rotor-furl loads, the new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Tail} = {}^E \mathbf{v}_r^W \cdot \mathbf{F}_{Tail}^W + {}^E \boldsymbol{\omega}_r^N \cdot \mathbf{M}_{Tail}^{N@W} \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with the tail and tail fin. Thus,

$$F_r|_{Tail} = F_r^*|_A + F_r|_{GravA} + F_r|_{AeroA} + F_r|_{SpringTF} + F_r|_{DampTF} \quad (r = 1, 2, \dots, 22)$$

Since ${}^E \mathbf{v}_r^W$ and ${}^E \boldsymbol{\omega}_r^N$ are equal to zero unless $r = 1, 2, \dots, 11$, the generalized active forces associated with tail-furl elasticity and damping do not contribute to the tail-furl loads (since also, $F_r|_{SpringTF}$ and $F_r|_{DampTF}$ are equal to zero if $r = 1, 2, \dots, 11$). So,

$$F_r|_{Tail} = F_r^*|_A + F_r|_{GravA} + F_r|_{AeroA} \quad (r = 1, 2, \dots, 11)$$

Thus,

$$F_r|_{Tail} = -m^B {}^E \mathbf{v}_r^I \cdot ({}^E \mathbf{a}^I + g\mathbf{z}_2) - m^F {}^E \mathbf{v}_r^J \cdot ({}^E \mathbf{a}^J + g\mathbf{z}_2) + {}^E \mathbf{v}_r^K \cdot \mathbf{F}_{AeroA}^K + {}^E \boldsymbol{\omega}_r^A \cdot (\mathbf{M}_{AeroA}^A - \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\alpha}^A + {}^E \boldsymbol{\omega}^A \times \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}^A) \quad (r = 1, 2, \dots, 11)$$

However, ${}^E \boldsymbol{\omega}_r^A$ and ${}^E \boldsymbol{\omega}_r^N$ are all equal when r is constrained to be between 1 and 11. Recognizing also that ${}^E \mathbf{v}_r^I = {}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WI}$, ${}^E \mathbf{v}_r^J = {}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WJ}$, and ${}^E \mathbf{v}_r^K = {}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WK}$, when $r = 1, 2, \dots, 11$, this generalized force can be expanded to:

$$F_r|_{Tail} = ({}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WK}) \cdot \mathbf{F}_{AeroA}^K - m^B ({}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WI}) \cdot ({}^E \mathbf{a}^I + g\mathbf{z}_2) - m^F ({}^E \mathbf{v}_r^W + {}^E \boldsymbol{\omega}_r^N \times \mathbf{r}^{WJ}) \cdot ({}^E \mathbf{a}^J + g\mathbf{z}_2) + {}^E \boldsymbol{\omega}_r^N \cdot (\mathbf{M}_{AeroA}^A - \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\alpha}^A - {}^E \boldsymbol{\omega}^A \times \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}^A) \quad (r = 1, 2, \dots, 11)$$

Now applying the cyclic permutation law of the scalar triple product:

$$F_r|_{Tail} = {}^E \mathbf{v}_r^W \cdot \left[\mathbf{F}_{AeroA}^K - m^B \left({}^E \mathbf{a}^I + g\mathbf{z}_2 \right) - m^F \left({}^E \mathbf{a}^J + g\mathbf{z}_2 \right) \right] \\ + {}^E \boldsymbol{\omega}_r^N \cdot \left(\mathbf{M}_{AeroA}^A + \mathbf{r}^{WK} \times \mathbf{F}_{AeroA}^K - m^B \mathbf{r}^{WI} \times \left({}^E \mathbf{a}^I + g\mathbf{z}_2 \right) - m^F \mathbf{r}^{WJ} \times \left({}^E \mathbf{a}^J + g\mathbf{z}_2 \right) - \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\alpha}^A - {}^E \boldsymbol{\omega}^A \times \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\omega}^A \right) \quad (r = 1, 2, \dots, 11)$$

Thus it is seen that,

$$\mathbf{F}_{Tail}^W = \mathbf{F}_{AeroA}^K - m^B \left({}^E \mathbf{a}^I + g\mathbf{z}_2 \right) - m^F \left({}^E \mathbf{a}^J + g\mathbf{z}_2 \right)$$

and

$$\mathbf{M}_{Tail}^{N@W} = \mathbf{M}_{AeroA}^A + \mathbf{r}^{WK} \times \mathbf{F}_{AeroA}^K - m^B \mathbf{r}^{WI} \times \left({}^E \mathbf{a}^I + g\mathbf{z}_2 \right) - m^F \mathbf{r}^{WJ} \times \left({}^E \mathbf{a}^J + g\mathbf{z}_2 \right) - \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\alpha}^A - {}^E \boldsymbol{\omega}^A \times \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\omega}^A$$

Thus,

$$\mathbf{F}_{Tail}^W = \mathbf{F}_{AeroA}^K - m^B \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^I \ddot{q}_i \right) + {}^E \mathbf{v}_{TFrl}^I \ddot{q}_{TFrl} + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^I \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{TFrl}^I \right) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} \\ - m^F \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^J \ddot{q}_i \right) + {}^E \mathbf{v}_{TFrl}^J \ddot{q}_{TFrl} + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^J \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{TFrl}^J \right) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Tail}^{N@W} = \mathbf{M}_{AeroA}^A + \mathbf{r}^{WK} \times \mathbf{F}_{AeroA}^K - m^B \mathbf{r}^{WI} \times \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^I \ddot{q}_i \right) + {}^E \mathbf{v}_{TFrl}^I \ddot{q}_{TFrl} + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^I \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{TFrl}^I \right) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} \\ - m^F \mathbf{r}^{WJ} \times \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^J \ddot{q}_i \right) + {}^E \mathbf{v}_{TFrl}^J \ddot{q}_{TFrl} + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^J \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \mathbf{v}_{TFrl}^J \right) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} \\ - \bar{\bar{\mathbf{I}}}^A \cdot \left\{ \left(\sum_{i=4}^{11} {}^E \boldsymbol{\omega}_i^A \ddot{q}_i \right) + {}^E \boldsymbol{\omega}_{TFrl}^A \ddot{q}_{TFrl} + \left[\sum_{i=7}^{11} \frac{d}{dt} \left({}^E \boldsymbol{\omega}_i^A \right) \dot{q}_i \right] + \frac{d}{dt} \left({}^E \boldsymbol{\omega}_{TFrl}^A \right) \dot{q}_{TFrl} \right\} - {}^E \boldsymbol{\omega}^A \times \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\omega}^A$$

Or,

$$\mathbf{F}_{Tail_r}^W = -m^B {}^E \mathbf{v}_r^I - m^F {}^E \mathbf{v}_r^J \quad (r = 1, 2, \dots, 11; 15)$$

$$\mathbf{F}_{Tail_i}^W = \mathbf{F}_{Aero}^K - m^B \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} ({}^E \mathbf{v}_i^I) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{TFrl}^I) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} - m^F \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} ({}^E \mathbf{v}_i^J) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{TFrl}^J) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Tail_r}^{N@W} = -m^B \mathbf{r}^{WI} \times {}^E \mathbf{v}_r^I - m^F \mathbf{r}^{WJ} \times {}^E \mathbf{v}_r^J - \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}_r^A \quad (r = 1, 2, \dots, 11; 15)$$

$$\begin{aligned} \mathbf{M}_{Tail_i}^{N@W} = & \mathbf{M}_{AeroA}^A + \mathbf{r}^{WK} \times \mathbf{F}_{AeroA}^K - m^B \mathbf{r}^{WI} \times \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} ({}^E \mathbf{v}_i^I) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{TFrl}^I) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} - m^F \mathbf{r}^{WJ} \times \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} ({}^E \mathbf{v}_i^J) \dot{q}_i \right] + \frac{d}{dt} ({}^E \mathbf{v}_{TFrl}^J) \dot{q}_{TFrl} + g\mathbf{z}_2 \right\} \\ & - \bar{\mathbf{I}}^A \cdot \left\{ \left[\sum_{i=7}^{11} \frac{d}{dt} ({}^E \boldsymbol{\omega}_i^A) \dot{q}_i \right] + \frac{d}{dt} ({}^E \boldsymbol{\omega}_{TFrl}^A) \dot{q}_{TFrl} \right\} - {}^E \boldsymbol{\omega}^A \times \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}^A \end{aligned}$$

The output loads is as follows,

$$TFrlBrM = \mathbf{M}_{Tail}^{N@W} \cdot \mathbf{tfa} / 1,000$$

Tail-furl bearing furl moment (about the tail-furl axis), (kN·m)

Like the *LSShftTq*, *LSSTipMza*, and *RFrlBrM*, it is noted that the tail-furling furl moment can be computed differently using the tail-furl springs and dampers, though the load summation method and this other constraint method are equivalent. This can be demonstrated as follows. First of all, the equation above is equivalent to saying:

$$TFrlBrM = {}^E \boldsymbol{\omega}_{TFrl}^A \cdot \mathbf{M}_{Tail}^{N@W} / 1,000$$

Or,

$$TFrlBrM = {}^E \boldsymbol{\omega}_{TFrl}^A \cdot \left[\mathbf{M}_{AeroA}^A + \mathbf{r}^{WK} \times \mathbf{F}_{AeroA}^K - m^B \mathbf{r}^{WI} \times ({}^E \mathbf{a}^I + g\mathbf{z}_2) - m^F \mathbf{r}^{WJ} \times ({}^E \mathbf{a}^J + g\mathbf{z}_2) - \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\alpha}^A - {}^E \boldsymbol{\omega}^A \times \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}^A \right] / 1,000$$

Now applying the cyclic permutation law of the scalar triple product:

$$TFrlBrM = \left[{}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WK} \cdot \mathbf{F}_{AeroA}^K - m^B {}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WI} \cdot ({}^E \mathbf{a}^I + g\mathbf{z}_2) - m^F {}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WJ} \cdot ({}^E \mathbf{a}^J + g\mathbf{z}_2) \right] / 1,000 \\ + {}^E \boldsymbol{\omega}_{TFrl}^A \cdot \left(\mathbf{M}_{AeroA}^A - \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\alpha}^A - {}^E \boldsymbol{\omega}^A \times \bar{\mathbf{I}}^A \cdot {}^E \boldsymbol{\omega}^A \right)$$

Recognizing also that ${}^E \mathbf{v}_{TFrl}^I = {}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WI}$, ${}^E \mathbf{v}_{TFrl}^J = {}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WJ}$, and ${}^E \mathbf{v}_{TFrl}^K = {}^E \boldsymbol{\omega}_{TFrl}^A \times \mathbf{r}^{WK}$, this can be expanded as follows:

$$TFrlBrM = \left[{}^E \mathbf{v}_{TFrl}^K \cdot \mathbf{F}_{AeroA}^K + {}^E \boldsymbol{\omega}_{TFrl}^A \cdot \mathbf{M}_{AeroA}^A - m^B {}^E \mathbf{v}_{TFrl}^I \cdot ({}^E \mathbf{a}^I + g\mathbf{z}_2) - m^F {}^E \mathbf{v}_{TFrl}^J \cdot ({}^E \mathbf{a}^J + g\mathbf{z}_2) - {}^E \boldsymbol{\omega}_{TFrl}^A \cdot (\bar{\bar{\mathbf{I}}}^A \cdot {}^E \mathbf{a}^A + {}^E \boldsymbol{\omega}^A \times \bar{\bar{\mathbf{I}}}^A \cdot {}^E \boldsymbol{\omega}^A) \right] / 1,000$$

or,

$$TFrlBrM = \left(F_{TFrl}^* \Big|_A + F_{TFrl} \Big|_{GravA} + F_{TFrl} \Big|_{AeroA} \right) / 1,000$$

From the equations of motion, it is easily seen that this is equivalent to saying:

$$TFrlBrM = \left(-F_{TFrl} \Big|_{SpringTF} - F_{TFrl} \Big|_{DampTF} \right) / 1,000$$

and thus,

$$TFrlBrM = \left\{ \begin{array}{l} TFrlSpr \cdot q_{TFrl} + IF \left[q_{TFrl} > TFrlUSSP, TFrlUSSpr(q_{TFrl} - TFrlUSSP), 0 \right] \\ + IF \left[q_{TFrl} < TFrlDSSP, TFrlDSSpr(q_{TFrl} - TFrlDSSP), 0 \right] \\ + TFrlDmp \cdot \dot{q}_{TFrl} + IF \left[\dot{q}_{TFrl} < 0, TFrlCDmp \cdot SIGN(\dot{q}_{TFrl}), 0 \right] \\ + IF \left[q_{TFrl} > TFrlUSDP, TFrlUSDmp \cdot \dot{q}_{TFrl}, 0 \right] + IF \left[q_{TFrl} < TFrlDSDP, TFrlDSDmp \cdot \dot{q}_{TFrl}, 0 \right] \end{array} \right\} / 1,000 \quad (= \mathbf{M}_{Tail}^{N@W} \cdot \mathbf{tfa} / 1,000)$$

Thus, both the load summation method and the constraint method are equivalent. Thus, to avoid using 2 different methods to calculate $TFrlBrM$ if

various DOFs are disabled, it is best just to use $\mathbf{M}_{Tail}^{N@W} \cdot \mathbf{tfa} / 1,000$, which will always work.

Tower Top / Yaw Bearing Loads:

There are 10 output loads at the tower top / yaw bearing location. 5 of them are the 3 components of tower top force $\mathbf{F}_{Nac,Rot}^O$ (2 components are expressed in a nonrotating frame, 2 components are expressed in a rotating frame, and 1 component is independent of rotation). The 5 other loads are the 3 components of the tower top bending moment, $\mathbf{M}_{Nac,Rot}^{B@O}$ (again, 2 components are expressed in a nonrotating frame, 2 components are expressed in a rotating frame, and 1 component is independent of rotation). All these loads are given relative to point O as indicated. Note that none of these loads include the effects of the yaw bearing mass (YawBrMass), which would affect the forces but not the moments. The new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Nac,Rot} = {}^E \mathbf{v}_r^O \cdot \mathbf{F}_{Nac,Rot}^O + {}^E \boldsymbol{\omega}_r^B \cdot \mathbf{M}_{Nac,Rot}^{B@O} \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with everything but the tower and platform. Thus,

$$\begin{aligned} F_r|_{Nac,Rot} = & F_r^*|_N + F_r^*|_R + F_r^*|_G + F_r^*|_H + F_r^*|_{B1} + F_r^*|_{B2} + F_r^*|_A \\ & + F_r|_{AeroB1} + F_r|_{AeroB2} + F_r|_{AeroA} + F_r|_{GravN} + F_r|_{GravR} + F_r|_{GravH} + F_r|_{GravB1} + F_r|_{GravB2} + F_r|_{GravA} + F_r|_{Gen} + F_r|_{Brake} + F_r|_{GBFric} \\ & + F_r|_{SpringYaw} + F_r|_{DampYaw} + F_r|_{SpringRF} + F_r|_{DampRF} + F_r|_{SpringTeet} + F_r|_{DampTeet} + F_r|_{SpringTF} + F_r|_{DampTF} \\ & + F_r|_{ElasticB1} + F_r|_{DampB1} + F_r|_{ElasticB2} + F_r|_{DampB2} + F_r|_{ElasticDrive} + F_r|_{DampDrive} \end{aligned} \quad (r = 1, 2, \dots, 22)$$

Since ${}^E \mathbf{v}_r^O$ and ${}^E \boldsymbol{\omega}_r^B$ are equal to zero unless $r = 1, 2, \dots, 10$, the generalized active forces associated with blade, drivetrain, yaw, rotor-furl, tail-furl, and teeter elasticity and damping as well as the generator torque, high-speed shaft braking torque, and gearbox friction do not contribute to the tower top loads (since also, $F_r|_{ElasticB1}$, $F_r|_{DampB1}$, $F_r|_{ElasticB2}$, $F_r|_{DampB2}$, $F_r|_{SpringTeet}$, $F_r|_{DampTeet}$, $F_r|_{SpringRF}$, $F_r|_{DampRF}$, $F_r|_{SpringTF}$, $F_r|_{DampTF}$, $F_r|_{SpringYaw}$, $F_r|_{DampYaw}$, $F_r|_{ElasticDrive}$, $F_r|_{DampDrive}$, $F_r|_{Gen}$, $F_r|_{Brake}$, and $F_r|_{GBFric}$ are equal to zero if $r = 1, 2, \dots, 10$). So,

$$\begin{aligned} F_r|_{Nac,Rot} = & F_r^*|_{B1} + F_r^*|_{AeroB1} + F_r^*|_{GravB1} + F_r^*|_{B2} + F_r^*|_{AeroB2} + F_r^*|_{GravB2} + F_r^*|_H + F_r^*|_{GravH} + F_r^*|_R + F_r^*|_{GravR} + F_r^*|_G \\ & + F_r^*|_A + F_r^*|_{GravA} + F_r^*|_{AeroA} + F_r^*|_N + F_r^*|_{GravN} \end{aligned} \quad (r = 1, 2, \dots, 10)$$

When using the results for the rotor-furl and tail-furl loads, this equation can be simplified as follows:

$$F_r|_{Nac,Rot} = F_r|_{Gen,Rot} + F_r|_{Tail} + F_r^*|_N + F_r|_{GravN} \quad (r = 1, 2, \dots, 10)$$

Thus,

$$F_r|_{Nac,Rot} = {}^E \mathbf{v}_r^V \cdot \mathbf{F}_{Gen,Rot}^V + {}^E \boldsymbol{\omega}_r^N \cdot \mathbf{M}_{Gen,Rot}^{N@V} + {}^E \mathbf{v}_r^W \cdot \mathbf{F}_{Tail}^W + {}^E \boldsymbol{\omega}_r^N \cdot \mathbf{M}_{Tail}^{N@W} - m^N {}^E \mathbf{v}_r^U \cdot ({}^E \mathbf{a}^U + g\mathbf{z}_2) - {}^E \boldsymbol{\omega}_r^N \cdot (\bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\alpha}^N + {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N) \quad (r = 1, 2, \dots, 10)$$

However, ${}^E \boldsymbol{\omega}_r^N$ and ${}^E \boldsymbol{\omega}_r^B$ are all equal when r is constrained to be between 1 and 10. Thus, when grouping like terms:

$$F_r|_{Nac,Rot} = {}^E \mathbf{v}_r^V \cdot \mathbf{F}_{Gen,Rot}^V + {}^E \mathbf{v}_r^W \cdot \mathbf{F}_{Tail}^W - m^N {}^E \mathbf{v}_r^U \cdot ({}^E \mathbf{a}^U + g\mathbf{z}_2) + {}^E \boldsymbol{\omega}_r^B \cdot (\mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} - \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\alpha}^N - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N) \quad (r = 1, 2, \dots, 10)$$

Recognizing also that ${}^E \mathbf{v}_r^U = {}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OU}$, ${}^E \mathbf{v}_r^V = {}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OV}$, and ${}^E \mathbf{v}_r^W = {}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OW}$, when $r = 1, 2, \dots, 10$, this generalized force can be expanded to:

$$F_r|_{Nac,Rot} = ({}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OV}) \cdot \mathbf{F}_{Gen,Rot}^V + ({}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OW}) \cdot \mathbf{F}_{Tail}^W - m^N ({}^E \mathbf{v}_r^O + {}^E \boldsymbol{\omega}_r^B \times \mathbf{r}^{OU}) \cdot ({}^E \mathbf{a}^U + g\mathbf{z}_2) + {}^E \boldsymbol{\omega}_r^B \cdot (\mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} - \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\alpha}^N - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N) \quad (r = 1, 2, \dots, 10)$$

Now applying the cyclic permutation law of the scalar triple product:

$$F_r|_{Nac,Rot} = {}^E \mathbf{v}_r^O \cdot [\mathbf{F}_{Gen,Rot}^V + \mathbf{F}_{Tail}^W - m^N ({}^E \mathbf{a}^U + g\mathbf{z}_2)] + {}^E \boldsymbol{\omega}_r^B \cdot [\mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail}^W - m^N \mathbf{r}^{OU} \times ({}^E \mathbf{a}^U + g\mathbf{z}_2)] + {}^E \boldsymbol{\omega}_r^B \cdot (\mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} - \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\alpha}^N - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N) \quad (r = 1, 2, \dots, 10)$$

which simplifies to:

$$F_r|_{Nac,Rot} = {}^E \mathbf{v}_r^O \cdot [\mathbf{F}_{Gen,Rot}^V + \mathbf{F}_{Tail}^W - m^N ({}^E \mathbf{a}^U + g\mathbf{z}_2)] + {}^E \boldsymbol{\omega}_r^B \cdot [\mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} + \mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail}^W - m^N \mathbf{r}^{OU} \times ({}^E \mathbf{a}^U + g\mathbf{z}_2) - \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\alpha}^N - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N] \quad (r = 1, 2, \dots, 10)$$

Thus it is seen that,

$$\mathbf{F}_{Nac,Rot}^O = \mathbf{F}_{Gen,Rot}^V + \mathbf{F}_{Tail}^W - m^N ({}^E \mathbf{a}^U + g\mathbf{z}_2)$$

and

$$\mathbf{M}_{Nac,Rot}^{B@O} = \mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} + \mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail}^W - m^N \mathbf{r}^{OU} \times \left({}^E \mathbf{a}^U + g\mathbf{z}_2 \right) - \bar{\mathbf{I}}^N \cdot {}^E \mathbf{a}^N - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N$$

Thus,

$$\mathbf{F}_{Nac,Rot}^O = \mathbf{F}_{Gen,Rot}^V + \mathbf{F}_{Tail}^W - m^N \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^U \ddot{q}_i \right) + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^U \right) \dot{q}_i \right] + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Nac,Rot}^{B@O} = \mathbf{M}_{Gen,Rot}^{N@V} + \mathbf{M}_{Tail}^{N@W} + \mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail}^W - m^N \mathbf{r}^{OU} \times \left\{ \left(\sum_{i=1}^{11} {}^E \mathbf{v}_i^U \ddot{q}_i \right) + \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^U \right) \dot{q}_i \right] + g\mathbf{z}_2 \right\} - \bar{\mathbf{I}}^N \cdot \left\{ \left(\sum_{i=4}^{11} {}^E \boldsymbol{\omega}_i^N \ddot{q}_i \right) + \left[\sum_{i=7}^{11} \frac{d}{dt} \left({}^E \boldsymbol{\omega}_i^N \right) \dot{q}_i \right] \right\} - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N$$

Or,

$$\mathbf{F}_{Nac,Rot_r}^O = \mathbf{F}_{Gen,Rot_r}^V + \mathbf{F}_{Tail_r}^W - m^N {}^E \mathbf{v}_r^U \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{F}_{Nac,Rot_i}^O = \mathbf{F}_{Gen,Rot_i}^V + \mathbf{F}_{Tail_i}^W - m^N \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^U \right) \dot{q}_i \right] + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{Nac,Rot_r}^{B@O} = \mathbf{M}_{Gen,Rot_r}^{N@V} + \mathbf{M}_{Tail_r}^{N@W} + \mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot_r}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail_r}^W - m^N \mathbf{r}^{OU} \times {}^E \mathbf{v}_r^U - \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}_r^N \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{M}_{Nac,Rot_i}^{B@O} = \mathbf{M}_{Gen,Rot_i}^{N@V} + \mathbf{M}_{Tail_i}^{N@W} + \mathbf{r}^{OV} \times \mathbf{F}_{Gen,Rot_i}^V + \mathbf{r}^{OW} \times \mathbf{F}_{Tail_i}^W - m^N \mathbf{r}^{OU} \times \left\{ \left[\sum_{i=4}^{11} \frac{d}{dt} \left({}^E \mathbf{v}_i^U \right) \dot{q}_i \right] + g\mathbf{z}_2 \right\} - \bar{\mathbf{I}}^N \cdot \left[\sum_{i=7}^{11} \frac{d}{dt} \left({}^E \boldsymbol{\omega}_i^N \right) \dot{q}_i \right] - {}^E \boldsymbol{\omega}^N \times \bar{\mathbf{I}}^N \cdot {}^E \boldsymbol{\omega}^N$$

The output loads are as follows,

$$\text{YawBrFxn} = \mathbf{F}_{Nac,Rot}^O \cdot \mathbf{d}_1 / 1,000$$

Rotating (with nacelle) yaw bearing shear force (directed along the xn-axis), (kN)

$$\text{YawBrFyn} = -\mathbf{F}_{Nac,Rot}^O \cdot \mathbf{d}_3 / 1,000$$

Rotating (with nacelle) yaw bearing shear force (directed along the yn-axis), (kN)

$$\text{YawBrFxp} = \mathbf{F}_{Nac,Rot}^O \cdot \mathbf{b}_1 / 1,000$$

Yaw bearing for-aft (nonrotating) shear force (directed along the xp-axis), (kN)

$$\text{YawBrFyp} = -\mathbf{F}_{Nac,Rot}^O \cdot \mathbf{b}_3 / 1,000$$

Yaw bearing side-to-side (nonrotating) shear force (directed along the yp-axis), (kN)

$$\text{YawBrFzn} = \text{YawBrFzp} = \mathbf{F}_{Nac,Rot}^O \cdot \mathbf{d}_2 / 1,000 = \mathbf{F}_{Nac,Rot}^O \cdot \mathbf{b}_2 / 1,000$$

Yaw bearing axial force (directed along the zn-/zp-axis), (kN)

$$\text{YawBrMxn} = \mathbf{M}_{Nac,Rot}^{B@O} \cdot \mathbf{d}_1 / 1,000$$

Rotating (with nacelle) yaw bearing roll moment (about the xn-axis), (kN·m)

$$YawBrMyn = -M_{Nac,Rot}^{B@O} \cdot d_3 / 1,000$$

Rotating (with nacelle) yaw bearing pitch moment (about the yn-axis), (kN·m)

$$YawBrMxp = M_{Nac,Rot}^{B@O} \cdot b_1 / 1,000$$

Nonrotating yaw bearing roll moment (about the xp-axis), (kN·m)

$$YawBrMyp = -M_{Nac,Rot}^{B@O} \cdot b_3 / 1,000$$

Nonrotating yaw bearing pitch moment (about the yp-axis), (kN·m)

$$YawBrMzn = YawBrMzp = M_{Nac,Rot}^{B@O} \cdot d_2 / 1,000 = M_{Nac,Rot}^{B@O} \cdot b_2 / 1,000$$

Yaw bearing yaw moment (about the zn-/zp-axis), (kN·m)

Like the *LSShftTq*, *LSSTipMza*, *RFrBrM*, and *TFrBrM*, it is noted that the yaw bearing yaw moment can be computed differently using the yaw drive spring and damper, though the load summation method and this other constraint method are equivalent. This can be demonstrated as follows. First of all, the equation above is equivalent to saying:

$$YawBrMzn = {}^E\omega_{Yaw}^N \cdot M_{Nac,Rot}^{B@O} / 1,000$$

Or,

$$YawBrMzn = {}^E\omega_{Yaw}^N \cdot \left[M_{Gen,Rot}^{N@V} + M_{Tail}^{N@W} + r^{OV} \times F_{Gen,Rot}^V + r^{OW} \times F_{Tail}^W - m^N r^{OU} \times ({}^Ea^U + gz_2) - \bar{I}^N \cdot {}^Ea^N - {}^E\omega^N \times \bar{I}^N \cdot {}^E\omega^N \right] / 1,000$$

Now applying the cyclic permutation law of the scalar triple product:

$$YawBrMzn = \left[{}^E\omega_{Yaw}^N \times r^{OV} \cdot F_{Gen,Rot}^V + {}^E\omega_{Yaw}^N \times r^{OW} \cdot F_{Tail}^W - m^N {}^E\omega_{Yaw}^N \times r^{OU} \cdot ({}^Ea^U + gz_2) + {}^E\omega_{Yaw}^N \cdot (M_{Gen,Rot}^{N@V} + M_{Tail}^{N@W} - \bar{I}^N \cdot {}^Ea^N - {}^E\omega^N \times \bar{I}^N \cdot {}^E\omega^N) \right] / 1,000$$

Recognizing also that ${}^E v_{Yaw}^U = {}^E\omega_{Yaw}^N \times r^{OU}$, ${}^E v_{Yaw}^V = {}^E\omega_{Yaw}^N \times r^{OV}$, and ${}^E v_{Yaw}^W = {}^E\omega_{Yaw}^N \times r^{OW}$, this can be expanded as follows:

$$YawBrMzn = \left[{}^E v_{Yaw}^V \cdot F_{Gen,Rot}^V + {}^E\omega_{Yaw}^N \cdot M_{Gen,Rot}^{N@V} + {}^E v_{Yaw}^W \cdot F_{Tail}^W + {}^E\omega_{Yaw}^N \cdot M_{Tail}^{N@W} - m^N {}^E v_{Yaw}^U \cdot ({}^Ea^U + gz_2) - {}^E\omega_{Yaw}^N \cdot (\bar{I}^N \cdot {}^Ea^N + {}^E\omega^N \times \bar{I}^N \cdot {}^E\omega^N) \right] / 1,000$$

or,

$$YawBrMzn = (F_{Yaw}|_{Rotor} + F_{Yaw}|_{Tail} + F_{Yaw}^*|_N + F_{Yaw}|_{GravN}) / 1,000$$

or,

$$YawBrMzn = \left(F_{Yaw}^*|_N + F_{Yaw}^*|_R + F_{Yaw}^*|_G + F_{Yaw}^*|_H + F_{Yaw}^*|_{B1} + F_{Yaw}^*|_{B2} + F_{Yaw}^*|_A + F_{Yaw}|_{AeroB1} + F_{Yaw}|_{AeroB2} + F_{Yaw}|_{AeroA} + F_{Yaw}|_{GravN} + F_{Yaw}|_{GravR} + F_{Yaw}|_{GravH} + F_{Yaw}|_{GravB1} + F_{Yaw}|_{GravB2} + F_{Yaw}|_{GravA} \right) / 1,000$$

From the equations of motion, it is easily seen that this is equivalent to saying:

$$YawBrMzn = \left(-F_{Yaw}|_{SpringYaw} - F_{Yaw}|_{DampYaw} \right) / 1,000$$

and thus,

$$YawBrMzn = \left[YawSpr(q_{Yaw} - YawNeut) + YawDamp \cdot \dot{q}_{Yaw} \right] / 1,000 \quad (= \mathbf{M}_{Nac,Rot}^{B@O} \cdot \mathbf{d}_2 / 1,000)$$

Thus, both the load summation method and the constraint method are equivalent. Thus, to avoid using 2 different methods to calculate $YawBrMzn$ if various DOFs are disabled, it is best just to use $\mathbf{M}_{Nac,Rot}^{B@O} \cdot \mathbf{d}_2 / 1,000$, which will always work.

Tower Base Loads:

There are 6 output loads at the base of the tower. 3 of them are the 3 components of the base force $\mathbf{F}_{Turb}^T(0)$. The other 3 are the 3 components of the base bending moments, $\mathbf{M}_{Turb}^X(0)$. Note that the tower base loads are all output at the point on the tower where it changes from being rigid to being flexible ($h = 0$). The new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{Turb} = {}^E \mathbf{v}_r^T(0) \cdot \mathbf{F}_{Turb}^T(0) + {}^E \boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Turb}^X(0) \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with everything but the platform. Thus,

$$\begin{aligned} F_r|_{Turb} = & F_r^*|_T + F_r^*|_N + F_r^*|_R + F_r^*|_G + F_r^*|_H + F_r^*|_{B1} + F_r^*|_{B2} + F_r^*|_A \\ & + F_r|_{AeroT} + F_r|_{AeroB1} + F_r|_{AeroB2} + F_r|_{AeroA} + F_r|_{GravT} + F_r|_{GravN} + F_r|_{GravR} + F_r|_{GravH} + F_r|_{GravB1} + F_r|_{GravB2} + F_r|_{GravA} \\ & + F_r|_{SpringYaw} + F_r|_{DampYaw} + F_r|_{SpringRF} + F_r|_{DampRF} + F_r|_{SpringTeet} + F_r|_{DampTeet} + F_r|_{SpringTF} + F_r|_{DampTF} + F_r|_{Gen} + F_r|_{Brake} + F_r|_{GBFric} \\ & + F_r|_{ElasticT} + F_r|_{DampT} + F_r|_{ElasticB1} + F_r|_{DampB1} + F_r|_{ElasticB2} + F_r|_{DampB2} + F_r|_{ElasticDrive} + F_r|_{DampDrive} \end{aligned} \quad (r = 1, 2, \dots, 22)$$

Since ${}^E \mathbf{v}_r^T(0)$ and ${}^E \boldsymbol{\omega}_r^X$ are equal to zero unless $r = 1, 2, \dots, 6$, the generalized active forces associated with blade, drivetrain, yaw, rotor-furl, tail-furl, teeter, and tower elasticity and damping as well as the generator torque, high-speed shaft braking torque, and gearbox friction do not contribute to the tower base loads (since also, $F_r|_{ElasticB1}, F_r|_{DampB1}, F_r|_{ElasticB2}, F_r|_{DampB2}, F_r|_{SpringTeet}, F_r|_{DampTeet}, F_r|_{SpringRF}, F_r|_{DampRF}, F_r|_{SpringTF}, F_r|_{DampTF}, F_r|_{SpringYaw}, F_r|_{DampYaw}, F_r|_{ElasticDrive}, F_r|_{DampDrive}, F_r|_{Gen}, F_r|_{Brake}, F_r|_{GBFric}, F_r|_{ElasticT}$, and $F_r|_{DampT}$ are equal to zero if $r = 1, 2, \dots, 6$). So,

$$\begin{aligned} F_r|_{Turb} = & F_r^*|_{B1} + F_r|_{AeroB1} + F_r|_{GravB1} + F_r^*|_{B2} + F_r|_{AeroB2} + F_r|_{GravB2} + F_r^*|_H + F_r|_{GravH} + F_r^*|_R + F_r|_{GravR} + F_r^*|_G \\ & + F_r^*|_A + F_r|_{GravA} + F_r|_{AeroA} + F_r^*|_N + F_r|_{GravN} + F_r^*|_T + F_r|_{AeroT} + F_r|_{GravT} \end{aligned} \quad (r = 1, 2, \dots, 6)$$

When using the results for the tower-top loads, this equation can be simplified as follows:

$$F_r|_{Turb} = F_r|_{Nac, Rot} + F_r^*|_T + F_r|_{AeroT} + F_r|_{GravT} \quad (r = 1, 2, \dots, 6)$$

Thus,

$$F_r|_{Turb} = {}^E \mathbf{v}_r^O \cdot \mathbf{F}_{Nac,Rot}^O + {}^E \boldsymbol{\omega}_r^B \cdot \mathbf{M}_{Nac,Rot}^{B@O} - \int_0^{TwrFlexL} \mu^T(h) {}^E \mathbf{v}_r^T(h) \cdot [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] dh - YawBrMass {}^E \mathbf{v}_r^O \cdot ({}^E \mathbf{a}^O + g\mathbf{z}_2) \\ + \int_0^{TwrFlexL} [{}^E \mathbf{v}_r^T(h) \cdot \mathbf{F}_{AeroT}^T(h) + {}^E \boldsymbol{\omega}_r^F(h) \cdot \mathbf{M}_{AeroT}^F(h)] dh \quad (r = 1, 2, \dots, 6)$$

However, ${}^E \boldsymbol{\omega}_r^B$, ${}^E \boldsymbol{\omega}_r^F(h)$, and ${}^E \boldsymbol{\omega}_r^X$ are all equal when r is constrained to be between 1 and 6. Thus,

$$F_r|_{Turb} = {}^E \mathbf{v}_r^O \cdot \mathbf{F}_{Nac,Rot}^O + {}^E \boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Nac,Rot}^{B@O} - \int_0^{TwrFlexL} \mu^T(h) {}^E \mathbf{v}_r^T(h) \cdot [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] dh - YawBrMass {}^E \mathbf{v}_r^O \cdot ({}^E \mathbf{a}^O + g\mathbf{z}_2) \\ + \int_0^{TwrFlexL} [{}^E \mathbf{v}_r^T(h) \cdot \mathbf{F}_{AeroT}^T(h) + {}^E \boldsymbol{\omega}_r^X(h) \cdot \mathbf{M}_{AeroT}^F(h)] dh \quad (r = 1, 2, \dots, 6)$$

Recognizing also that ${}^E \mathbf{v}_r^O = {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)]$, and ${}^E \mathbf{v}_r^T(h) = {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)]$, when $r = 1, 2, \dots, 6$, this generalized force can be expanded to:

$$F_r|_{Turb} = \left\{ {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \right\} \cdot \mathbf{F}_{Nac,Rot}^O + {}^E \boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Nac,Rot}^{B@O} \\ - \int_0^{TwrFlexL} \mu^T(h) \left\{ {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \right\} \cdot [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] dh \\ - YawBrMass \left\{ {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \right\} \cdot ({}^E \mathbf{a}^O + g\mathbf{z}_2) \\ + \int_0^{TwrFlexL} \left(\left\{ {}^E \mathbf{v}_r^T(0) + {}^E \boldsymbol{\omega}_r^X \times [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \right\} \cdot \mathbf{F}_{AeroT}^T(h) + {}^E \boldsymbol{\omega}_r^X(h) \cdot \mathbf{M}_{AeroT}^F(h) \right) dh \quad (r = 1, 2, \dots, 6)$$

Now applying the cyclic permutation law of the scalar triple product and simplifying:

$$F_r|_{Turb} = {}^E \mathbf{v}_r^T(0) \cdot \left(\mathbf{F}_{Nac,Rot}^O - YawBrMass({}^E \mathbf{a}^O + g\mathbf{z}_2) + \int_0^{TwrFlexL} \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] \right\} dh \right) \\ + {}^E \boldsymbol{\omega}_r^X \cdot \left(\mathbf{M}_{Nac,Rot}^{B@O} + [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \times [\mathbf{F}_{Nac,Rot}^O - YawBrMass({}^E \mathbf{a}^O + g\mathbf{z}_2)] \right. \\ \left. + \int_0^{TwrFlexL} [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \times \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] \right\} dh + \int_0^{TwrFlexL} \mathbf{M}_{AeroT}^F(h) dh \right) \quad (r = 1, 2, \dots, 6)$$

Thus it is seen that,

$$\mathbf{F}_{Turb}^T(0) = \mathbf{F}_{Nac,Rot}^O - YawBrMass({}^E \mathbf{a}^O + g\mathbf{z}_2) + \int_0^{TwrFlexL} \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] \right\} dh$$

and

$$\mathbf{M}_{Turb}^X(0) = \mathbf{M}_{Nac,Rot}^{B@O} + [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \times [\mathbf{F}_{Nac,Rot}^O - YawBrMass({}^E \mathbf{a}^O + g\mathbf{z}_2)] \\ + \int_0^{TwrFlexL} [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \times \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) [{}^E \mathbf{a}^T(h) + g\mathbf{z}_2] \right\} dh + \int_0^{TwrFlexL} \mathbf{M}_{AeroT}^F(h) dh$$

Thus,

$$\mathbf{F}_{Turb}^T(0) = \mathbf{F}_{Nac,Rot}^O - YawBrMass \left\{ \left(\sum_{i=1}^{10} {}^E \mathbf{v}_i^O \ddot{q}_i \right) + \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^O) \dot{q}_i \right] + g\mathbf{z}_2 \right\} + \int_0^{TwrFlexL} \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) \left\{ \left(\sum_{i=1}^{10} {}^E \mathbf{v}_i^T(h) \ddot{q}_i \right) + \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^T(h)) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right\} dh$$

and

$$\mathbf{M}_{Turb}^X(0) = \mathbf{M}_{Nac,Rot}^{B@O} + [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \times \left(\mathbf{F}_{Nac,Rot}^O - YawBrMass \left\{ \left(\sum_{i=1}^{10} {}^E \mathbf{v}_i^O \ddot{q}_i \right) + \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^O) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right) \\ + \int_0^{TwrFlexL} [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \times \left(\mathbf{F}_{AeroT}^T(h) - \mu^T(h) \left\{ \left(\sum_{i=1}^{10} {}^E \mathbf{v}_i^T(h) \ddot{q}_i \right) + \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^T(h)) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right) dh + \int_0^{TwrFlexL} \mathbf{M}_{AeroT}^F(h) dh$$

Or,

$$\mathbf{F}_{Turb_r}^T(0) = \mathbf{F}_{Nac, Rot_r}^O - YawBrMass {}^E \mathbf{v}_r^O - \int_0^{TwrFlexL} \mu^T(h) {}^E \mathbf{v}_r^T(h) dh \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{F}_{Turb_i}^T(0) = \mathbf{F}_{Nac, Rot_i}^O - YawBrMass \left\{ \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^O) \dot{q}_i \right] + g\mathbf{z}_2 \right\} + \int_0^{TwrFlexL} \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) \left\{ \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^T(h)) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right\} dh$$

and

$$\mathbf{M}_{Turb_r}^X(0) = \mathbf{M}_{Nac, Rot_r}^{B@O} + [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \times (\mathbf{F}_{Nac, Rot_r}^O - YawBrMass {}^E \mathbf{v}_r^O) - \int_0^{TwrFlexL} [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \times [\mu^T(h) {}^E \mathbf{v}_r^T(h)] dh \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{M}_{Turb_i}^X(0) = \mathbf{M}_{Nac, Rot_i}^{B@O} + [\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(0)] \times \left(\mathbf{F}_{Nac, Rot_i}^O - YawBrMass \left\{ \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^O) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right)$$

$$+ \int_0^{TwrFlexL} [\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(0)] \times \left(\mathbf{F}_{AeroT}^T(h) - \mu^T(h) \left\{ \left[\sum_{i=4}^{10} \frac{d}{dt} ({}^E \mathbf{v}_i^T(h)) \dot{q}_i \right] + g\mathbf{z}_2 \right\} \right) dh + \int_0^{TwrFlexL} \mathbf{M}_{AeroT}^F(h) dh$$

Thus,

$$TwrBsFxt = \mathbf{F}_{Turb}^T(0) \cdot \mathbf{a}_1 / 1,000$$

Tower base fore-aft shear force (directed along the xt-axis), (kN)

$$TwrBsFyt = -\mathbf{F}_{Turb}^T(0) \cdot \mathbf{a}_3 / 1,000$$

Tower base side-to-side shear force (directed along the yt-axis), (kN)

$$TwrBsFzt = \mathbf{F}_{Turb}^T(0) \cdot \mathbf{a}_2 / 1,000$$

Tower base axial force (directed along the zt-axis), (kN)

$$TwrBsMxt = \mathbf{M}_{Turb}^X(0) \cdot \mathbf{a}_1 / 1,000$$

Tower base roll (or side-to-side) moment (i.e., the moment caused by side-to-side forces) (about the xt-axis),

(kN·m)

$$TwrBsMyt = -\mathbf{M}_{Turb}^X(0) \cdot \mathbf{a}_3 / 1,000$$

Tower base pitching (or fore-aft) moment (i.e., the moment caused by fore-aft forces) (about the yt-axis), (kN·m)

$$TwrBsMzt = \mathbf{M}_{Turb}^X(0) \cdot \mathbf{a}_2 / 1,000$$

Tower base yaw (or torsional) moment (about the zt-axis), (kN·m)

Tower Local Moment Outputs:

There are 3 output loads at any of the selected tower node locations i ($h = H^{Node\ i}$) ($i=1,2,\dots,5$). These are the 3 components of the bending moment $\mathbf{M}_{Turb}^F(H^{Node\ i})$ expressed in the *local* tower element coordinate system (principal structural axes). Examining the results for the tower base loads, it follows that:

$$\mathbf{M}_{Turb}^F(H^{Node\ i}) = \mathbf{M}_{Nac, Rot}^{B@O} + \left[\mathbf{r}^{ZO} - \mathbf{r}^{ZT}(H^{Node\ i}) \right] \times \left[\mathbf{F}_{Nac, Rot}^O - YawBrMass \left({}^E \mathbf{a}^O + g\mathbf{z}_2 \right) \right] \\ + \int_0^{TwrFlexL} \left[\mathbf{r}^{ZT}(h) - \mathbf{r}^{ZT}(H^{Node\ i}) \right] \times \left\{ \mathbf{F}_{AeroT}^T(h) - \mu^T(h) \left[{}^E \mathbf{a}^T(h) + g\mathbf{z}_2 \right] \right\} dh + \int_{H^{Node\ i}}^{TwrFlexL} \mathbf{M}_{AeroT}^F(h) dh \quad (i = 1, 2, \dots, 5)$$

The output loads are as follows:

$TwHtiMLxt = \mathbf{M}_{Turb}^F(H^{Node\ i}) \cdot \mathbf{t}_1^{B1}(H^{Node\ i}) / 1,000$	Tower <i>local</i> roll moment of tower gage i (about the <i>local</i> xt-structural axis), (kN·m)
$TwHtiMLyt = -\mathbf{M}_{Turb}^F(H^{Node\ i}) \cdot \mathbf{t}_3^{B1}(H^{Node\ i}) / 1,000$	Tower <i>local</i> pitching moment of tower gage i (about the <i>local</i> yt-structural axis), (kN·m)
$TwHtiMLzt = \mathbf{M}_{Turb}^F(H^{Node\ i}) \cdot \mathbf{t}_2^{B1}(H^{Node\ i}) / 1,000$	Tower <i>local</i> yaw (or torsion) moment of tower gage i (about the <i>local</i> zt-structural axis), (kN·m)

Platform Loads:

There are 12 output loads at the platform reference point. 6 of them are the 3 components of the platform force $\mathbf{F}_{Hydro}^Z$ (3 components expressed in the tower base / platform coordinate system and 3 components expressed in the inertia frame). The remaining 6 are the 3 components of the platform moment $\mathbf{M}_{Hydro}^{X@Z}$ (3 components expressed in the tower base / platform coordinate system and 3 components expressed in the inertia frame). These are the loads transmitted from the water/mooring lines or foundation to the platform.

The output loads are as follows:

$$PtfmFxt = \mathbf{F}_{Hydro}^Z \cdot \mathbf{a}_1 / 1,000$$

Platform horizontal surge force (directed along the xt-axis), (kN)

$$PtfmFyt = -\mathbf{F}_{Hydro}^Z \cdot \mathbf{a}_3 / 1,000$$

Platform horizontal sway force (directed along the yt-axis), (kN)

$$PtfmFzt = \mathbf{F}_{Hydro}^Z \cdot \mathbf{a}_2 / 1,000$$

Platform vertical heave force (directed along the zt-axis), (kN)

$$PtfmFxi = \mathbf{F}_{Hydro}^Z \cdot \mathbf{z}_1 / 1,000$$

Platform horizontal surge force (directed along the xi-axis), (kN)

$$PtfmFyi = -\mathbf{F}_{Hydro}^Z \cdot \mathbf{z}_3 / 1,000$$

Platform horizontal sway force (directed along the yi-axis), (kN)

$$PtfmFzi = \mathbf{F}_{Hydro}^Z \cdot \mathbf{z}_2 / 1,000$$

Platform vertical heave force (directed along the zi-axis), (kN)

$$PtfmMxt = \mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{a}_1 / 1,000$$

Platform roll tilt moment (about the xt-axis), (kN·m)

$$PtfmMyt = -\mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{a}_3 / 1,000$$

Platform pitch tilt moment (about the yt-axis), (kN·m)

$$PtfmMzt = \mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{a}_2 / 1,000$$

Platform yaw moment (about the zt-axis), (kN·m)

$$PtfmMxi = \mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{z}_1 / 1,000$$

Platform roll tilt moment (about the xi-axis), (kN·m)

$$PtfmMyi = -\mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{z}_3 / 1,000$$

Platform pitch tilt moment (about the yi-axis), (kN·m)

$$PtfmMzi = \mathbf{M}_{Hydro}^{X@Z} \cdot \mathbf{z}_2 / 1,000$$

Platform yaw moment (about the zi-axis), (kN·m)

However, there are two loads, $\mathbf{F}_{All}^Z$ and $\mathbf{M}_{All}^{X@Z}$, that are useful to use when assembling the equations of motion. *Both of these loads are always equal zero*, defining the balance between all inertia loads and all applied forces. It makes the most sense to discuss these loads in this section. The new generalized active force for the equations of motion resulting from these new loads is:

$$F_r|_{All} = {}^E \mathbf{v}_r^Z \cdot \mathbf{F}_{All}^Z + {}^E \boldsymbol{\omega}_r^X \cdot \mathbf{M}_{All}^{X@Z} \quad (r = 1, 2, \dots, 22)$$

This generalized active force must produce the same effects as the generalized active and inertia forces associated with *everything*. Since ${}^E\mathbf{v}_r^Z$ and ${}^E\boldsymbol{\omega}_r^X$ are equal to zero unless $r = 1, 2, \dots, 6$, nothing but inertia, gravity, aerodynamics, and hydrodynamics contribute to these loads. So,

$$F_r|_{All} = F_r^*|_{Bl} + F_r|_{AeroBl} + F_r|_{GravBl} + F_r^*|_{B2} + F_r|_{AeroB2} + F_r|_{GravB2} + F_r^*|_H + F_r|_{GravH} + F_r^*|_R + F_r|_{GravR} + F_r^*|_G + F_r^*|_A + F_r|_{GravA} + F_r|_{AeroA} + F_r^*|_N + F_r|_{GravN} + F_r^*|_T + F_r|_{AeroT} + F_r|_{GravT} + F_r^*|_X + F_r|_{GravX} + F_r|_{HydroX} \quad (r = 1, 2, \dots, 6)$$

When using the results for the tower base loads, this equation can be simplified as follows:

$$F_r|_{All} = F_r|_{Turb} + F_r^*|_X + F_r|_{GravX} + F_r|_{HydroX} \quad (r = 1, 2, \dots, 6)$$

Thus,

$$F_r|_{All} = {}^E\mathbf{v}_r^T(0) \cdot \mathbf{F}_{Turb}^T(0) + {}^E\boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Turb}^X(0) + {}^E\mathbf{v}_r^Z \cdot \mathbf{F}_{Hydro}^Z + {}^E\boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Hydro}^{X@Z} - m^X {}^E\mathbf{v}_r^Y \cdot ({}^E\mathbf{a}^Y + g\mathbf{z}_2) - {}^E\boldsymbol{\omega}_r^X \cdot (\bar{\bar{\mathbf{I}}}^X \cdot {}^E\mathbf{a}^X + {}^E\boldsymbol{\omega}_r^X \times \bar{\bar{\mathbf{I}}}^X \cdot {}^E\boldsymbol{\omega}_r^X) \quad (r = 1, 2, \dots, 6)$$

Recognizing also that ${}^E\mathbf{v}_r^T(0) = {}^E\mathbf{v}_r^Z + {}^E\boldsymbol{\omega}_r^X \times \mathbf{r}^{ZT}(0)$ and ${}^E\mathbf{v}_r^Y = {}^E\mathbf{v}_r^Z + {}^E\boldsymbol{\omega}_r^X \times \mathbf{r}^{ZY}$, when $r = 1, 2, \dots, 6$, this generalized force can be expanded to:

$$F_r|_{All} = \left[{}^E\mathbf{v}_r^Z + {}^E\boldsymbol{\omega}_r^X \times \mathbf{r}^{ZT}(0) \right] \cdot \mathbf{F}_{Turb}^T(0) + {}^E\boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Turb}^X(0) + {}^E\mathbf{v}_r^Z \cdot \mathbf{F}_{Hydro}^Z + {}^E\boldsymbol{\omega}_r^X \cdot \mathbf{M}_{Hydro}^{X@Z} - m^X ({}^E\mathbf{v}_r^Z + {}^E\boldsymbol{\omega}_r^X \times \mathbf{r}^{ZY}) \cdot ({}^E\mathbf{a}^Y + g\mathbf{z}_2) - {}^E\boldsymbol{\omega}_r^X \cdot (\bar{\bar{\mathbf{I}}}^X \cdot {}^E\mathbf{a}^X + {}^E\boldsymbol{\omega}_r^X \times \bar{\bar{\mathbf{I}}}^X \cdot {}^E\boldsymbol{\omega}_r^X) \quad (r = 1, 2, \dots, 6)$$

Now applying the cyclic permutation law of the scalar triple product and simplifying:

$$F_r|_{All} = {}^E\mathbf{v}_r^Z \cdot \left[\mathbf{F}_{Turb}^T(0) + \mathbf{F}_{Hydro}^Z - m^X ({}^E\mathbf{a}^Y + g\mathbf{z}_2) \right] + {}^E\boldsymbol{\omega}_r^X \cdot \left[\mathbf{M}_{Turb}^X(0) + \mathbf{M}_{Hydro}^{X@Z} + \mathbf{r}^{ZT}(0) \times \mathbf{F}_{Turb}^T(0) - m^X \mathbf{r}^{ZY} \times ({}^E\mathbf{a}^Y + g\mathbf{z}_2) - \bar{\bar{\mathbf{I}}}^X \cdot {}^E\mathbf{a}^X - {}^E\boldsymbol{\omega}_r^X \times \bar{\bar{\mathbf{I}}}^X \cdot {}^E\boldsymbol{\omega}_r^X \right] \quad (r = 1, 2, \dots, 6)$$

Thus it is seen that,

$$\mathbf{F}_{All}^Z = \mathbf{F}_{Turb}^T(0) + \mathbf{F}_{Hydro}^Z - m^X \left({}^E \mathbf{a}^Y + g\mathbf{z}_2 \right)$$

and

$$\mathbf{M}_{All}^{X@Z} = \mathbf{M}_{Turb}^X(0) + \mathbf{M}_{Hydro}^{X@Z} + \mathbf{r}^{ZT}(0) \times \mathbf{F}_{Turb}^T(0) - m^X \mathbf{r}^{ZY} \times \left({}^E \mathbf{a}^Y + g\mathbf{z}_2 \right) - \bar{\mathbf{I}}^X \cdot {}^E \boldsymbol{\alpha}^X - {}^E \boldsymbol{\omega}^X \times \bar{\mathbf{I}}^X \cdot {}^E \boldsymbol{\omega}^X$$

Thus,

$$\mathbf{F}_{All}^Z = \mathbf{F}_{Turb}^T(0) + \left(\sum_{j=1}^6 \mathbf{F}_{Hydro_j}^Z \ddot{q}_j \right) + \mathbf{F}_{Hydro_t}^Z - m^X \left\{ \left(\sum_{i=1}^6 {}^E \mathbf{v}_i^Y \ddot{q}_i \right) + \left[\sum_{i=4}^6 \frac{d}{dt} ({}^E \mathbf{v}_i^Y) \dot{q}_i \right] + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{All}^{X@Z} = \mathbf{M}_{Turb}^X(0) + \left(\sum_{j=1}^6 \mathbf{M}_{Hydro_j}^{X@Z} \ddot{q}_j \right) + \mathbf{M}_{Hydro_t}^{X@Z} + \mathbf{r}^{ZT}(0) \times \mathbf{F}_{Turb}^T(0) - m^X \mathbf{r}^{ZY} \times \left\{ \left(\sum_{i=1}^6 {}^E \mathbf{v}_i^Y \ddot{q}_i \right) + \left[\sum_{i=4}^6 \frac{d}{dt} ({}^E \mathbf{v}_i^Y) \dot{q}_i \right] + g\mathbf{z}_2 \right\} - \bar{\mathbf{I}}^X \cdot \left(\sum_{i=4}^6 {}^E \boldsymbol{\omega}_i^X \ddot{q}_i \right) - {}^E \boldsymbol{\omega}^X \times \bar{\mathbf{I}}^X \cdot {}^E \boldsymbol{\omega}^X$$

Or,

$$\mathbf{F}_{All_r}^Z = \mathbf{F}_{Turb_r}^T(0) + \mathbf{F}_{Hydro_r}^Z - m^X \left[{}^E \mathbf{v}_r^Y \right] \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{F}_{All_r}^Z = \mathbf{F}_{Turb_r}^T(0) + \mathbf{F}_{Hydro_r}^Z - m^X \left\{ \left[\sum_{i=4}^6 \frac{d}{dt} ({}^E \mathbf{v}_i^Y) \dot{q}_i \right] + g\mathbf{z}_2 \right\}$$

and

$$\mathbf{M}_{All_r}^{X@Z} = \mathbf{M}_{Turb_r}^X(0) + \mathbf{r}^{ZT}(0) \times \mathbf{F}_{Turb_r}^T(0) + \mathbf{M}_{Hydro_r}^{X@Z} - m^X \mathbf{r}^{ZY} \times {}^E \mathbf{v}_r^Y - \bar{\mathbf{I}}^X \cdot {}^E \boldsymbol{\omega}_r^X \quad (r = 1, 2, \dots, 22)$$

$$\mathbf{M}_{All_t}^{X@Z} = \mathbf{M}_{Turb_t}^X(0) + \mathbf{r}^{ZT}(0) \times \mathbf{F}_{Turb_t}^T(0) + \mathbf{M}_{Hydro_t}^{X@Z} - m^X \mathbf{r}^{ZY} \times \left\{ \left[\sum_{i=4}^6 \frac{d}{dt} ({}^E \mathbf{v}_i^Y) \dot{q}_i \right] + g\mathbf{z}_2 \right\} - {}^E \boldsymbol{\omega}^X \times \bar{\mathbf{I}}^X \cdot {}^E \boldsymbol{\omega}^X$$

Equations of Motion In Terms of Loads:

The reason for finding the partial loads is that many portions of the equations of motion can be expressed in terms of the partial loads instead of in the form given in “FASTKinetics.doc”. Incorporating the partial loads into the development of the equations of motion is beneficial since it requires less computation time to compute the loads if members of the loads were already found when compiling the equations of motion. For example, many integrations must be made to develop the portions of the equations of motion associated with the blades. Several more integrations must be made in order to find the loads once the accelerations are found. When using partial loads associated with the blades to develop the equations of motion, the additional integrals to find the loads once the accelerations are found will be unnecessary. The main point is that the equations of motion and the output loads are inherently coupled, and the entire simulation can be done with fewer computations if the system of equations is developed with the load outputs in mind.

Examining the results from the previous sections of this document, it is easy to see that many portions of the equations of motion can be written in terms of the partial loads as follows:

$$\begin{aligned}
 & \left[[C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} \right] (Teet,r), (r,Teet) = -{}^E \omega_{Teet}^H \cdot M_{Rotor_r}^{L@P} \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22) \\
 & \left[[C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} \right] (DrTr,r), (r,DrTr) = -{}^E \omega_{DrTr}^L \cdot M_{Rotor_r}^{L@P} \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22) \\
 & \left[[C(q,t)]_R + [C(q,t)]_G + [C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} \right] (RFrl,r), (r,RFrl) = -{}^E \omega_{RFrl}^R \cdot M_{Gen,Rot_r}^{N@V} \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22) \\
 & \left[[C(q,t)]_N + [C(q,t)]_R + [C(q,t)]_G + [C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} + [C(q,t)]_A \right] (Yaw,r), (r,Yaw) = -{}^E \omega_{Yaw}^N \cdot M_{Nac,Rot_r}^{B@O} \quad (r = 1, 2, \dots, 22) \\
 & \left[[C(q,t)]_N + [C(q,t)]_R + [C(q,t)]_G + [C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} + [C(q,t)]_A \right] (a,r), (r,a) = -{}^E v_a^O \cdot F_{Nac,Rot_r}^O - {}^E \omega_a^B \cdot M_{Nac,Rot_r}^{B@O} \quad (a = 7, 8, \dots, 10, r = 1, 2, \dots, 22) \\
 & \left[[C(q,t)]_X + [C(q,t)]_{HydroX} + [C(q,t)]_T + [C(q,t)]_N + [C(q,t)]_R + [C(q,t)]_G + [C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} + [C(q,t)]_A \right] (a,r), (r,a) = -{}^E \omega_a^X \cdot M_{All_r}^{X@Z} \quad (a = 4, 5, 6, r = 1, 2, \dots, 22) \\
 & \left[[C(q,t)]_X + [C(q,t)]_{HydroX} + [C(q,t)]_T + [C(q,t)]_N + [C(q,t)]_R + [C(q,t)]_H + [C(q,t)]_{B1} + [C(q,t)]_{B2} + [C(q,t)]_A \right] (a,r), (r,a) = -{}^E v_a^Z \cdot F_{All_r}^Z \quad (a = 1, 2, 3, r = 1, 2, \dots, 22) \\
 & [C(q,t)]_A (TFrl,r), (r,TFrl) = -{}^E \omega_{TFrl}^A \cdot M_{Tail_r}^{N@W} \quad (r = 1, 2, \dots, 11; 15)
 \end{aligned}$$

← This last expression is not needed; instead, just add $\left[C(q,t) \right]_A (15,15)$.

Also, since DOFs DrTr and GeAz are so similar,

$$\begin{aligned} & \left[C(q,t) \right]_H + \left[C(q,t) \right]_{B1} + \left[C(q,t) \right]_{B2} \left(GeAz, r \right), (r, GeAz) \\ & = \left[C(q,t) \right]_H + \left[C(q,t) \right]_{B1} + \left[C(q,t) \right]_{B2} \left(DrTr, r \right), (r, DrTr) \end{aligned} \quad (r = 1, 2, \dots, 14; 16, 17, \dots, 22)$$

← This last expression is only used for $(r = 13, 14; 16, 17, \dots, 22)$ however. This is because if this expression was used for all of the r 's then the $\left[C(q,t) \right]_G$ effects for the generator azimuth DOF row and column would be removed for $(r = 4, 5, \dots, 12)$, which is undesirable.

The only additional terms that need to be added to the overall mass matrix are as follows:

$$\left[C(q,t) \right]_{B1} (Row, Col = 16, 17, 18), \left[C(q,t) \right]_{B2} (Row, Col = 19, 20, 21), \left[C(q,t) \right]_T (Row, Col = 7, 8, \dots, 10), \left[C(q,t) \right]_G (15, 15), \text{ and } \left[C(q,t) \right]_{GBFric}$$

Also,

$$\begin{aligned} & \left\{ \begin{aligned} & -f(\dot{q}, q, t) \Big|_H + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravH} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB1} \\ & + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB2} \end{aligned} \right\} \left(Teet \right) = {}^E \omega_{Teet}^H \cdot M_{Rotor_i}^{L@P} \\ & \left\{ \begin{aligned} & -f(\dot{q}, q, t) \Big|_H + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravH} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB1} \\ & + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB2} \end{aligned} \right\} \left(DrTr \right) = {}^E \omega_{DrTr}^L \cdot M_{Rotor_i}^{L@P} \\ & \left\{ \begin{aligned} & -f(\dot{q}, q, t) \Big|_R + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravR} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_G + \left\{ -f(\dot{q}, q, t) \right\} \Big|_H + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravH} \\ & + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB1} \\ & + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{B2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\} \Big|_{AeroB2} \end{aligned} \right\} \left(RFrI \right) = {}^E \omega_{RFrI}^R \cdot M_{Gen, Rot_i}^{N@V} \end{aligned}$$

$$\left\{ \begin{aligned}
 & \left\{ -f(\dot{q}, q, t) \right\}_N + \left\{ -f(\dot{q}, q, t) \right\}_{GravN} + \left\{ -f(\dot{q}, q, t) \right\}_R + \left\{ -f(\dot{q}, q, t) \right\}_{GravR} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_G + \left\{ -f(\dot{q}, q, t) \right\}_H + \left\{ -f(\dot{q}, q, t) \right\}_{GravH} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B1} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB1} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B2} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB2} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_A + \left\{ -f(\dot{q}, q, t) \right\}_{GravA} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroA}
 \end{aligned} \right\} (Yaw) = {}^E \omega_{Yaw}^N \cdot M_{Nac, Rot_t}^{B@O}$$

$$\left\{ \begin{aligned}
 & \left\{ -f(\dot{q}, q, t) \right\}_N + \left\{ -f(\dot{q}, q, t) \right\}_{GravN} + \left\{ -f(\dot{q}, q, t) \right\}_R + \left\{ -f(\dot{q}, q, t) \right\}_{GravR} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_G + \left\{ -f(\dot{q}, q, t) \right\}_H + \left\{ -f(\dot{q}, q, t) \right\}_{GravH} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B1} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB1} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B2} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB2} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_A + \left\{ -f(\dot{q}, q, t) \right\}_{GravA} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroA}
 \end{aligned} \right\} (Row) = {}^E v_{Row}^O \cdot F_{Nac, Rot_t}^O + {}^E \omega_{Row}^B \cdot M_{Nac, Rot_t}^{B@O} \quad (Row = 7, 8, \dots, 10)$$

$$\left\{ \begin{aligned}
 & \left\{ -f(\dot{q}, q, t) \right\}_X + \left\{ -f(\dot{q}, q, t) \right\}_{HydroX} + \left\{ -f(\dot{q}, q, t) \right\}_{GravX} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_T + \left\{ -f(\dot{q}, q, t) \right\}_{GravT} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroT}
 \end{aligned} \right\}$$

$$\left\{ \begin{aligned}
 & + \left\{ -f(\dot{q}, q, t) \right\}_N + \left\{ -f(\dot{q}, q, t) \right\}_{GravN} + \left\{ -f(\dot{q}, q, t) \right\}_R + \left\{ -f(\dot{q}, q, t) \right\}_{GravR} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_G + \left\{ -f(\dot{q}, q, t) \right\}_H + \left\{ -f(\dot{q}, q, t) \right\}_{GravH} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B1} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB1} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB1} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{B2} + \left\{ -f(\dot{q}, q, t) \right\}_{GravB2} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroB2} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_A + \left\{ -f(\dot{q}, q, t) \right\}_{GravA} + \left\{ -f(\dot{q}, q, t) \right\}_{AeroA}
 \end{aligned} \right\} (Row) = {}^E \omega_{Row}^X \cdot M_{All_t}^{X@Z} \quad (Row = 4, 5, 6)$$

$$\left\{ \begin{aligned}
 & \left\{ -f(\dot{q}, q, t) \right\}_{|_X} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{HydroX}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravX}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_T} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravT}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroT}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_N} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravN}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_R} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravR}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_G} + \left\{ -f(\dot{q}, q, t) \right\}_{|_H} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravH}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_{B1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB1}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_{B2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB2}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_A} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravA}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroA}}
 \end{aligned} \right\} (Row) = {}^E \mathbf{v}_{Row}^Z \cdot \mathbf{F}_{All}^Z \quad (Row = 1, 2, 3)$$

$$\left\{ \begin{aligned}
 & \left\{ -f(\dot{q}, q, t) \right\}_{|_H} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravH}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{B1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB1}} \\
 & + \left\{ -f(\dot{q}, q, t) \right\}_{|_{B2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB2}}
 \end{aligned} \right\} (GeAz) = {}^E \omega_{GeAz}^L \cdot \mathbf{M}_{Rotor}^{L@P}$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_A} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravA}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroA}} \} (TFrI) = {}^E \omega_{TFrI}^A \cdot \mathbf{M}_{Tail}^{N@W}$$

The only additional terms that need to be added to the overall forcing function are as follows:

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_{B1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB1}} \} (Row = 16, 17, 18), \left\{ -f(\dot{q}, q, t) \right\}_{|_{ElasticB1}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampB1}},$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_{B2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravB2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroB2}} \} (Row = 19, 20, 21), \left\{ -f(\dot{q}, q, t) \right\}_{|_{ElasticB2}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampB2}},$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_{SpringTeet}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampTeet}},$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_{SpringRF}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampRF}}, \left\{ -f(\dot{q}, q, t) \right\}_{|_{SpringTF}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampTF}}, \left\{ -f(\dot{q}, q, t) \right\}_{|_{SpringYaw}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampYaw}},$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_T} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{AeroT}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{GravT}} \} (Row = 7, 8, \dots, 10), \left\{ -f(\dot{q}, q, t) \right\}_{|_{ElasticT}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampT}},$$

$$\left\{ -f(\dot{q}, q, t) \right\}_{|_G} (13), \left\{ -f(\dot{q}, q, t) \right\}_{|_{Gen}}, \left\{ -f(\dot{q}, q, t) \right\}_{|_{Brake}}, \left\{ -f(\dot{q}, q, t) \right\}_{|_{GBFrict}}, \text{ and } \left\{ -f(\dot{q}, q, t) \right\}_{|_{ElasticDrive}} + \left\{ -f(\dot{q}, q, t) \right\}_{|_{DampDrive}}.$$